

# Shell evolution and collectivity in model spaces above $^{78}\text{Ni}$ and $^{132}\text{Sn}$ cores

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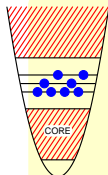
International Collaborations in Nuclear Theory:  
Theory for open-shell nuclei near the limits of stability

26.05.2015

# Shell model approach

## Calculations Ab Initio

- Realistic NN interactions
- Diagonalization in  $N\hbar\omega$  h.o.space



- define valence space
- $H_{\text{eff}}\Psi_{\text{eff}} = E\Psi_{\text{eff}}$
- ↪ INTERACTIONS
- build and diagonalize Hamiltonian matrix
- ↪ CODES

## Weak processes:

- $\beta$  decays
- $\beta\beta$  decays

$$[T_{1/2}^{0\nu}(0^+ \rightarrow 0^+)]^{-1} = G_{0\nu} |M^{0\nu}|^2 \langle m_\nu \rangle^2$$

☛ ASTROPHYSICS

☛ PARTICLE PHYSICS

## Collective excitations:

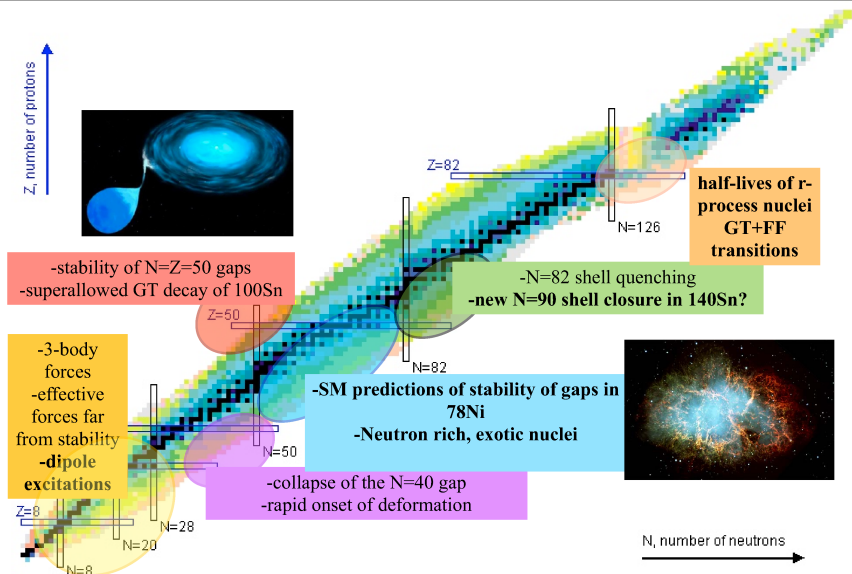
- deformation, superdeformation
- superfluidity
- symmetries

## Shell evolution far from stability:

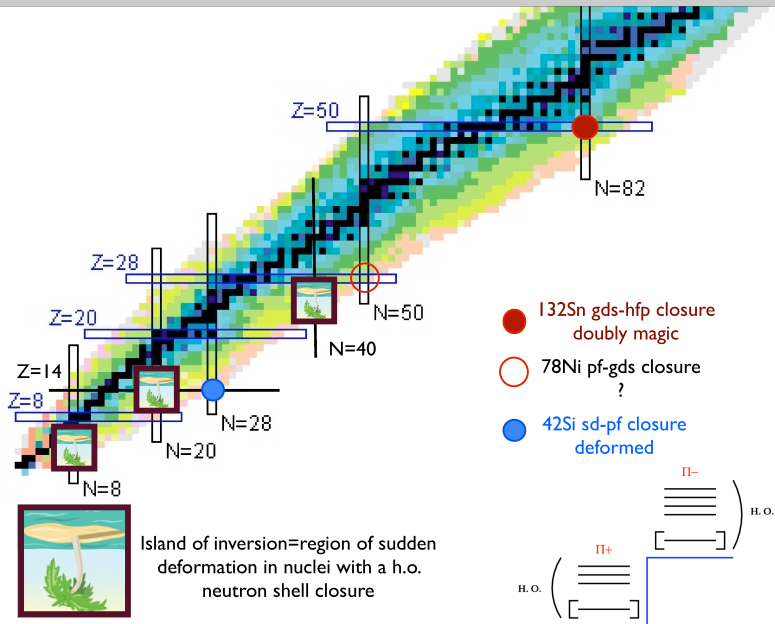
- Shell quenching
- New magic numbers

☛ ASTROPHYSICS

# SM with empirical interactions: regions of activity



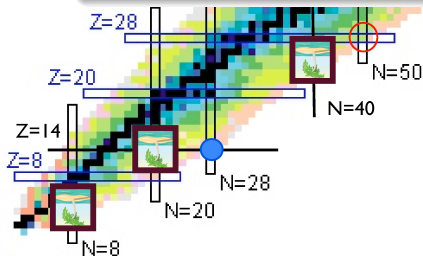
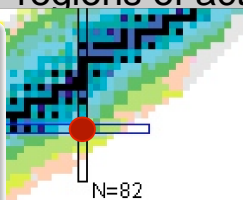
# Shell evolution: what have we learned ?



# SM with empirical interactions: regions of activity

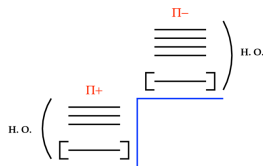
## PLAN:

- ① New results on copper isotopes & Z=28 gap
- ② Collectivity and triaxiality of nuclei above  $^{78}\text{Ni}$  and  $^{132}\text{Sn}$  cores
- ③ Some remarks about pairing in heavy tin region

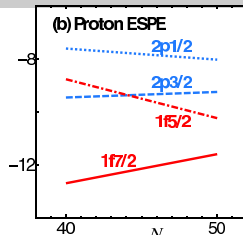
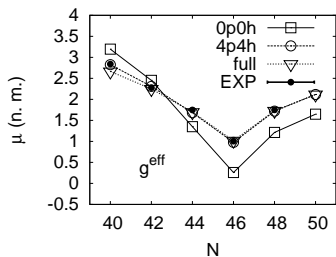


Island of inversion = region of sudden deformation in nuclei with a h.o. neutron shell closure

- $^{132}\text{Sn}$  gds-hfp closure doubly magic
- $^{78}\text{Ni}$  pf-gds closure ?
- $^{42}\text{Si}$  sd-pf closure deformed

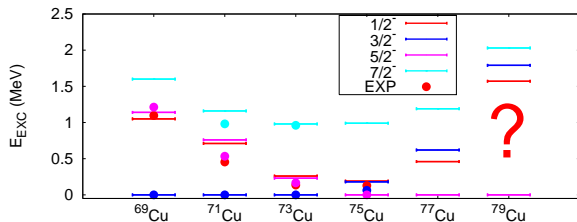


# Shell evolution towards $^{78}\text{Ni}$ : Z=28 gap

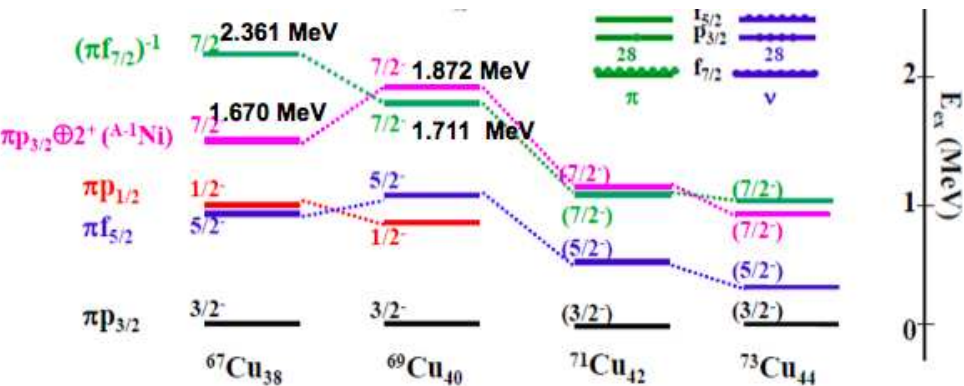


- crucial role of proton-core excitation for the description of magnetic moments

- crossing of  $\pi p_{3/2}-f_{5/2}$  orbitals with the filling of the  $\nu g_{9/2}$
- reduction of the Z = 28 gap in  $^{78}\text{Ni}$  due to the tensor force



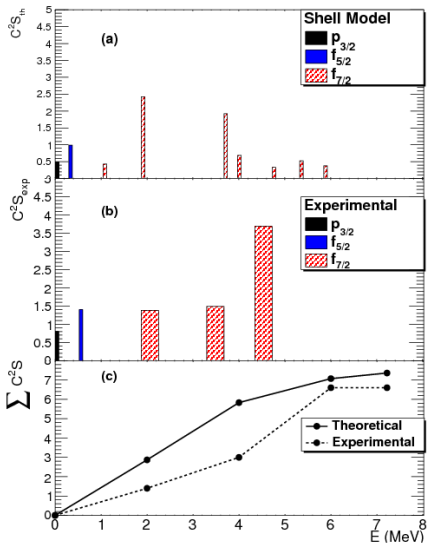
# Shell evolution in coppers: $7/2^-$ states



- Calculations in  $\pi pf$ - $\nu fpgd$  model space with updated LNPS interaction

*S. M. Lenzi et al., Phys. Rev. C82 (2010) 054301*

E (MeV)	$J^\pi$	Percentage	Composition
0	$3/2^-$	60%	$0_v^+ \otimes \pi p_{3/2}$
		14%	$2_v^+ \otimes \pi p_{3/2}$
0.31	$5/2^-$	36%	$0_v^+ \otimes \pi f_{5/2}$
		22%	$4_v^+ \otimes \pi f_{5/2}$
1.09	$7/2^-$	47%	$2_v^+ \otimes \pi p_{3/2}$
		11%	$4_v^+ \otimes \pi p_{3/2}$
		10%	$J_v^+ \otimes \pi f_{7/2}^{-1}$
1.41	$7/2^-$	42%	$2_v^+ \otimes \pi f_{5/2}$
		12%	$4_v^+ \otimes \pi f_{5/2}$
		13%	$J_v^+ \otimes \pi f_{7/2}^{-1}$



*P. Morfouace, S. Franchoo, K. Sieja et al., submitted*



# Shell evolution in copper isotopes: $^{69,71,73}\text{Cu}$ - B(E2)

Nucleus	$J_f^\pi \rightarrow J_i^\pi$	$E_\gamma$ (keV)	B(E2) ( $e^2\text{fm}^4$ )	EXP
$^{69}\text{Cu}$	$5/2_1^- \rightarrow 3/2_1^-$	1156	28.5	50(8)
	$7/2_1^- \rightarrow 3/2_1^-$	1804	46	77(12)
	$7/2_2^- \rightarrow 3/2_1^-$	2101	0.05	1.5
$^{71}\text{Cu}$	$5/2_1^- \rightarrow 3/2_1^-$	276	54	68(8.8)
	$7/2_1^- \rightarrow 3/2_1^-$	1115	179	187(2)
	$7/2_2^- \rightarrow 3/2_1^-$	1426	14.4	44(20)?
	$3/2_2^- \rightarrow 3/2_1^-$	1144	10.4	44(20)?
$^{73}\text{Cu}$	$5/2_1^- \rightarrow 3/2_1^-$	62	65	80(9)
	$7/2_1^- \rightarrow 3/2_1^-$	875	251	270(3)
	$7/2_2^- \rightarrow 5/2_1^-$	1128	0.02	20(11)?
	$5/2_2^- \rightarrow 5/2_1^-$	1203	11.4	20(11)?

$\pi p_{3/2} \otimes \nu 2^+$

$\pi f_{5/2} \otimes \nu 2^+$

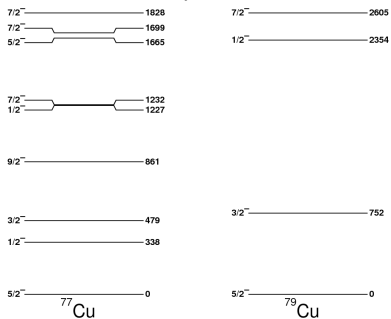
$\pi f_{7/2}^{-1}$

- SM and B(E2)'s consistent with non-observation of low lying strength in  $^{71}\text{Cu}$

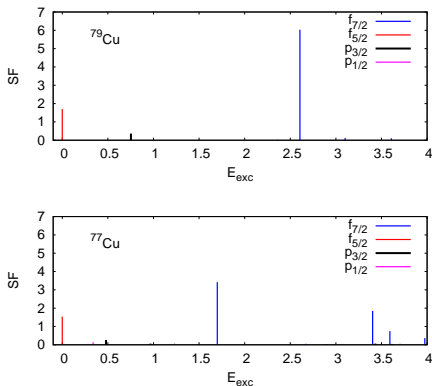
*E. Sahin, M. Doncel, K. Sieja et al., Phys. Rev. 91 (2015) 034302*

# Towards $^{78}\text{Ni}$ : $^{77,79}\text{Cu}$

Present SM: spectra

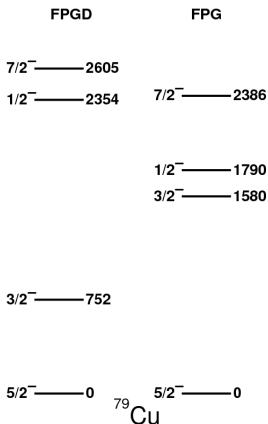
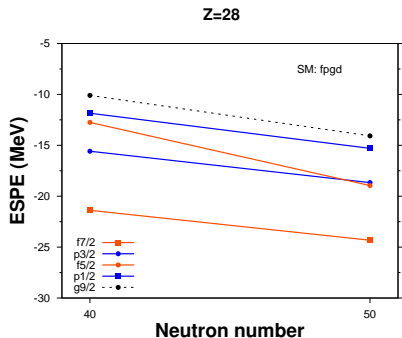


Present SM: SF distributions



● Data from knock-out reactions from RIKEN soon available

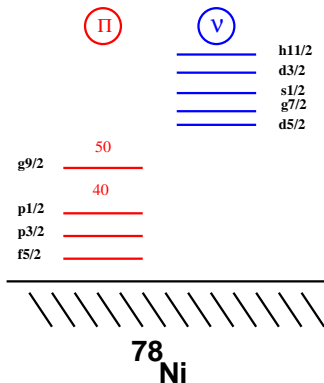
# Evolution of proton ESPE in copper isotopes



- Tensor effects seem to be exaggerated (or should be better compensated) in previous interaction in the *fpg* model space.
- 350keV impact on the predicted  $2^+$  energy of  $^{78}\text{Ni}$  (but anyhow missing neutron core excitations).

→ see the talk of F. Nowacki for the latest news ☺

# Nuclei above $^{78}\text{Ni}$

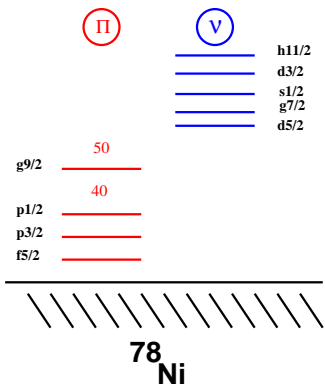


Knowing s.p. structure of  $^{79}\text{Ni}$  will be useful to validate the model assumptions

**Interaction:**  $\pi\pi$  fit of Lisetskiy & Brown,  $\nu\nu$  GCN5082,  $\pi\nu$  monopole corrected G-matrix. Proven successful and predictive in a large number of applications:

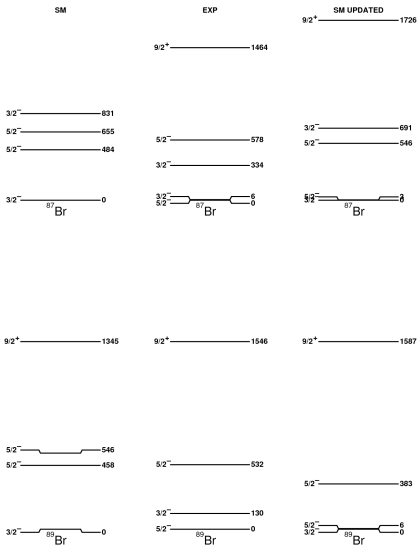
- Structure, mixed symmetry states in Zr isotopes, shell evolution between  $^{91}\text{Zr}$  and  $^{101}\text{Sn}$   
*K. Sieja et al., Phys. Rev. C79 (2009) 064310*
- Isomers and medium-spin structures of  $^{95}\text{Y}$ ,  $^{91-95}\text{Rb}$ ,  $^{92-96}\text{Sr}$   
*PRC85 (2012) 014329, PRC79 (2009) 024319, PRC82 (2010) 024302, PRC79 (2009) 044304*
- Collectivity and j-1 anomaly of  $^{87}\text{Se}$   
*PRC88 (2013) 034302*
- $\beta$ -decays of Ga nuclei and structure of  $N = 52, 54$  isotones  
*PRC88 (2013) 047301, PRC88 (2013) 044330, PRC88 (2013) 044314*
- Magnetic moments of  $^{86}\text{Kr}$ ,  $^{88}\text{Sr}$ ,  
*PRC 80 (2014) 064305*
- **Collectivity of  $N = 52, 54$  nuclei**  
*PRC88 (2013) 034327*

# Nuclei above $^{78}\text{Ni}$ : new developments



- New estimate of proton  $f_{5/2}$ - $p_{3/2}$  splitting in the core
- New fit of proton-proton interaction for N=50 isotones, using some new data e.g.  $^{83}\text{As}$  from EXILL

*P. Baczyk et al., to be published*

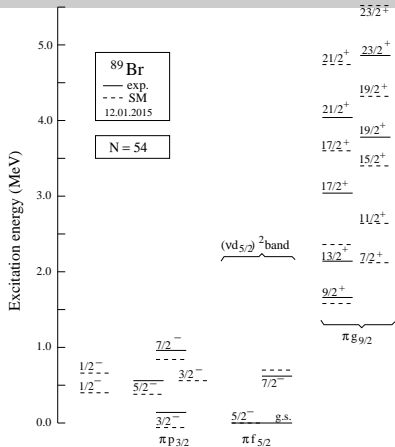
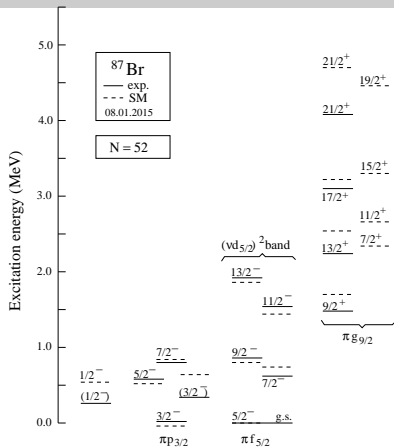


N=52

N=54

# Nuclei above $^{78}\text{Ni}$ : new developments

## Structure of odd and odd-odd nuclei



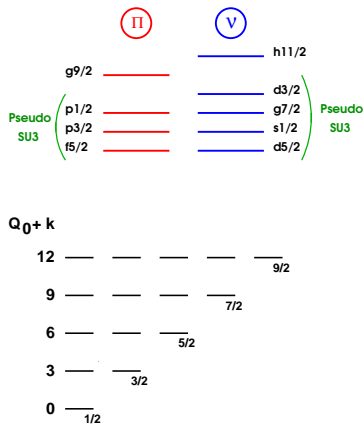
- Satisfactory ( $rms = 200\text{keV}$ ) description of proton-neutron multiplets in odd-odd nuclei:  $^{86,88}\text{Br}$ ,  $^{90}\text{Rb}$ .

*M. Czerwinski et al., to appear in Phys. Rev. C*

*M. Czerwinski et al., to be submitted.*

# Collectivity and triaxiality above $^{78}\text{Ni}$

Simple perspective: predictions of the pseudo-SU(3) scheme



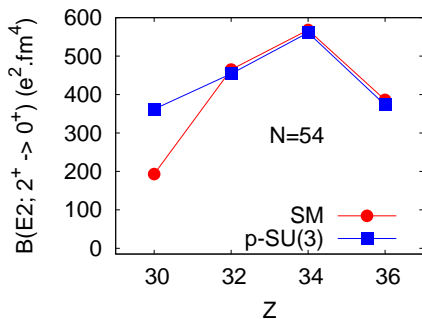
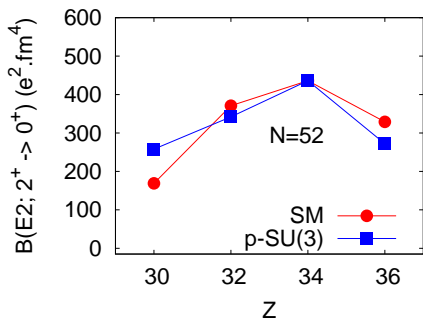
Predictions of pseudo-SU(3) model for  $N=52,54$ :

Nucleus	$Q_0$	$B(E2; 2^+ \rightarrow 0^+)$
$^{82}\text{Zn}$	114	258
$^{84}\text{Ge}$	131	342
$^{86}\text{Se}$	148	436
$^{88}\text{Kr}$	117	272
$^{84}\text{Zn}$	135	362
$^{86}\text{Ge}$	151	454
$^{88}\text{Se}$	168	561
$^{90}\text{Kr}$	137	373

- maximal prolate deformation in  $^{88}\text{Se}$
- possible triaxiality in  $^{86}\text{Ge}$

A. P. Zuker et al., Nilsson-SU3 selfconsistency: Quadrupole dominance in heavy  $N=Z$  nuclei, <http://arxiv.org/abs/1404.0224>

# Predictions of the pseudo-SU(3) model vs SM diagonalization



■ pseudo-SU(3) is a good approximation for the proton mid-shell

- maximal prolate deformation in  $^{88}\text{Se}$
- possible triaxiality in  $^{86}\text{Ge}$



# Intrinsic shape parameters of shell model states

(Yrast)	State	$Q_0$	$\beta$	$\gamma$ (deg.)
$^{86}\text{Ge}$	$0_{gs}^+$	165	0.238	12
	$2_1^+$	161	0.232	8
	$4_1^+$	152	0.218	12
	$6_1^+$	118	0.172	10
$^{88}\text{Se}$	$0_{gs}^+$	174	0.250	9
	$2_1^+$	169	0.243	12
	$4_1^+$	159	0.229	15
	$6_1^+$	118	0.173	14
(Excited)	State	$Q_0$	$\beta$	$\gamma$ (deg.)
$^{86}\text{Ge}$	$2_2^+$	152	0.219	28
	$3_1^+$	148	0.213	32
	$4_2^+$	116	0.169	41
	$5_1^+$	105	0.154	33
	$^{88}\text{Se}$	$2_2^+$	152	0.219
$3_2^+$		143	0.207	36
$4_2^+$		114	0.166	40
$5_2^+$		100	0.146	36

$$Q_{int}(s) = \sqrt{\frac{16\pi}{5}} p_s^{(2)}$$

$$\cos 3\gamma(s) = -\sqrt{7/2} p_s^{(3)} (p_s^{(2)})^{-3/2}$$

$$\begin{aligned} p_s^{(2)} &= (2I_s + 1)^{-1} \sum_r M_{sr}^2 \\ &= \frac{5(I_s + 1)(2I_s + 3)}{16\pi I_s(2I_s - 1)} Q_{spec}^2(s) \\ &+ \sum_{r \neq s} B(E2; s \rightarrow r), \end{aligned}$$

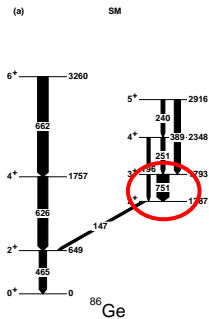
$M_{sr}$  = reduced E2 matrix elements

$$\begin{aligned} p_s^{(3)} &= -\sqrt{5}(2I_s + 1)^{-1} (-1)^{2I_s} \\ &\sum_{rt} \left\{ \begin{array}{ccc} 2 & 2 & 2 \\ I_s & I_r & I_t \end{array} \right\} M_{sr} M_{rt} M_{ts}, \end{aligned}$$

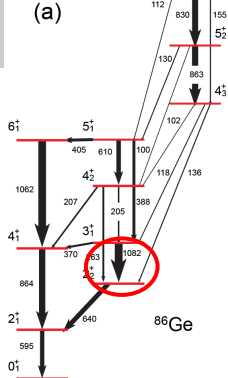
*K. Kumar, Phys. Rev. Lett. 28 (1972) 249*



# Triaxiality above $^{78}\text{Ni}$ ?



Shell Model using  $^{78}\text{Ni}$  core



GCM-Gogny

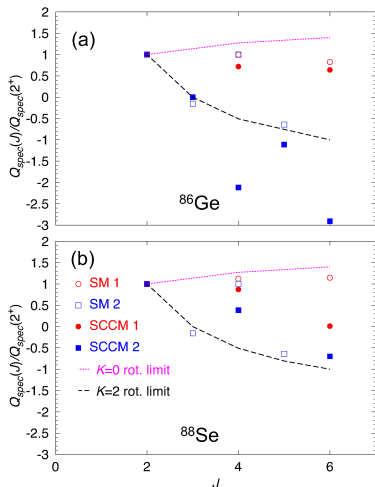
Agreement of excitation energies for the 1st excited band **within keV!**

- matrices dimension  $10^6$
- feasible on a laptop
- typical time of calculations:  
5min to 4h on one processor

- symmetry conserved (particle number and angular momentum)
- cluster of 140 CPUs
- typical time of calculations: 1 month

*K. Sieja, T.R. Rodriguez, K. Kolos and D. Verney, Phys. Rev. C88 (2013) 034327*

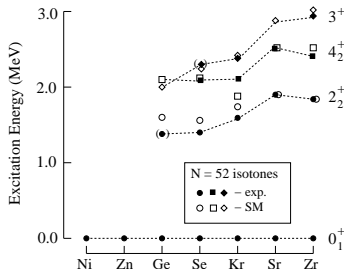
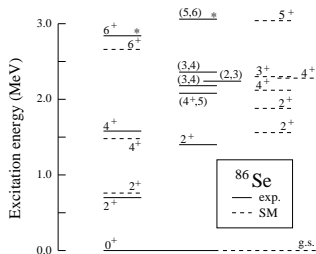
# Deformation at N=54 from different models



- $^{88}\text{Se}$  more spherical than  $^{86}\text{Ge}$  in SCCM
- the opposite in SM
- where is located the  $2^+$  of  $^{88}\text{Se}$ ? 886 keV or 651 keV ?
- both SCCM and SM agree  $^{86}\text{Ge}$  is the most triaxial
- both SCCM and SM predict  $\gamma$ -band does not continue to  $4_2^+$

# Some experimental results at N=52

T. Materna et al., to be published



- Second excited state in  $^{84}\text{Ge}$  is  $2_2^+$ , in agreement with predictions of 5DCH-D1S Gogny model, which gives a  $\gamma$ -soft ground state. SM supports the Gogny predictions with  $2_2^+$  being a head of a quasi- $\gamma$  band.

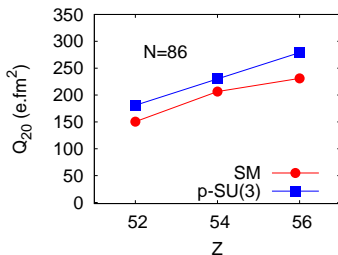
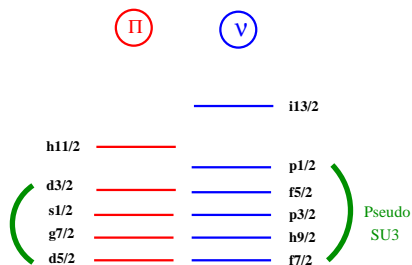
K. Kolos et al., *Phys. Rev. C*88 (2013) 047301; A. Korgul et al., *Phys. Rev. C*88 (2013) 044330

- $^{86}\text{Se}$  appears little deformed ( $\beta \sim 0.1$ ) but  $\gamma$ -soft from analysis of n-body quadrupole moments.
- Recent B(E2) from lifetimes (Cologne) in  $^{86}\text{Se}$  in good agreement with the SM ones!

# Collectivity above $^{132}\text{Sn}$ core

## pseudo-SU(3) predictions

In analogy to what we have seen above the  $^{78}\text{Ni}$ , collectivity should also thrive for the open-shell nuclei above the  $^{132}\text{Sn}$  core ( $N=54 \rightarrow N=86$ ). In the pseudo-SU(3) scheme (pseudo *pf* for protons and pseudo *gds* for neutrons) one obtains:

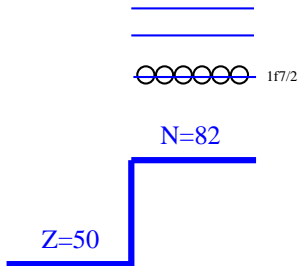


- Diagonalization in proton *gds* -neutron *hfp* model space, interaction under development.
- Some K-mixing should be possible in realistic calculations leading to triaxially deformed shapes.

● SM results in good agreement with experiment in  $^{138}\text{Te}$  and  $^{140}\text{Xe}$  (EXILL data) where an emerging  $\gamma$ -band was possibly observed.

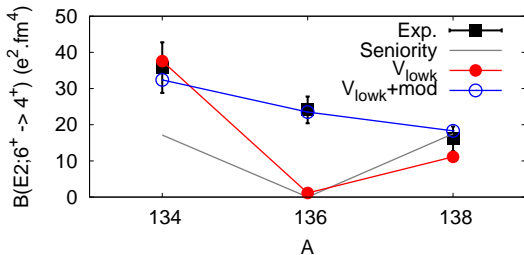
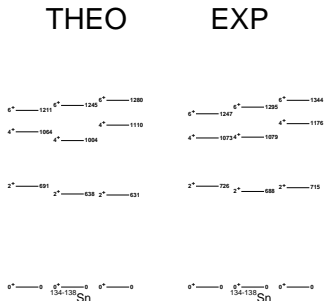
# Above $^{132}\text{Sn}$

- Taking  $^{132}\text{Sn}$  as a core (or  $^{110}\text{Zr}$ , see talk of H. Naidja) proton  $r4h$  and neutron  $r5i$  model space
- S.p. energies as usually used by other authors:
  - $\varepsilon_{p_{3/2}}=0.854$  MeV
  - $\varepsilon_{p_{1/2}}=1.363$  MeV
  - $\varepsilon_{f_{5/2}}=2.005$  MeV
  - $\varepsilon_{h_{9/2}}=1.561$  MeV
  - $\varepsilon_{i_{13/2}}=2.694$  MeV
- TBME derived from CD-Bonn (N3LO) potentials using  $V_{lowk}$  and MBPT ( $10\hbar\omega$  in  $\hat{Q}$ -box)



# Seniority isomers in the tin chain

Exp - PhD work of G. Gey (LPSC Grenoble)  
G. Simpson et al., PRL 113, 132502 (2014)



- Previous SM calculations wrong by a factor 2 for the  $B(E2; 6^+ \rightarrow 4^+)$  in  $^{136}\text{Sn}$
- reduction of  $f_{7/2}$  diagonal and non-diagonal pairing by 150 keV in the  $V_{lowk} + mod$  interaction

$$B(E2; \nu J \rightarrow \nu J - 2) \sim \left(1 - \frac{2n}{(2j+1)}\right)^2$$



# Isomeric transitions in tins: seniority mixing

Table: Dominating seniority components of the low lying states in tin isotopes

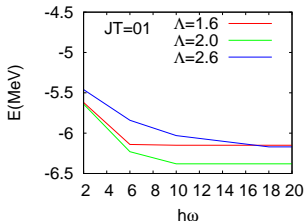
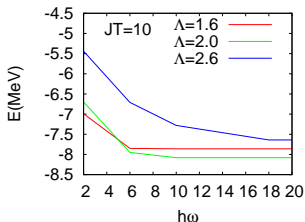
$J^\pi$	seniority	$^{136}\text{Sn}$	$^{138}\text{Sn}$
$0^+$	$\nu = 0$	98%	96%
	$\nu = 2$	2%	4%
$2^+$	$\nu = 2$	93%	87%
	$\nu = 4$	7%	10%
$4^+$	$\nu = 2$	46%	84%
	$\nu = 4$	54%	12%
$6^+$	$\nu = 2$	95%	87%
	$\nu = 4$	5%	9%

- Minor role of core excitations  
→ talk of H. Naidja
- The seniority mixing effects were observed important in other cases, e.g.  $^{72,74}\text{Ni}$

# Interactions from MBPT & pairing

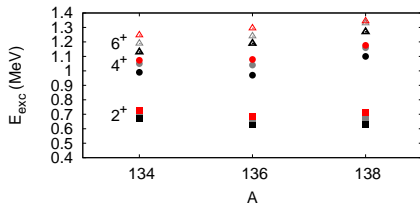
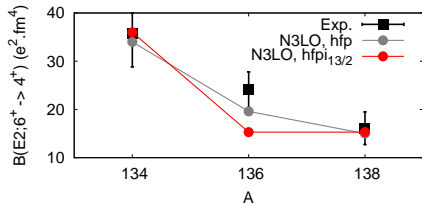
- *B. Maheshwari, A. K. Jain, P.C. Srivastava, Phys. Rev. C91 (2015) 024321*  
reduction of pairing matrix elements by 25keV (instead of our 150keV) in the RCDB interaction from OXBASH package to get the agreement on isomeric transitions
- our interaction includes 2nd order diagrams up to  $10\hbar\omega$  in  $\hat{Q}$ -box
- which interaction is better suited for tin isotopes?

Convergence of pairing eigenvalue in the  $sd$ -shell with number of  $\hbar\omega$  excitations in the  $\hat{Q}$ -box



# Interactions from MBPT & pairing

## Role of the intruder orbital



*N. Tsunoda et al., Multishell effective interactions, Phys. Rev. C89 (2014) 024313*

■ would be interesting to check the heavy tin region!

■ what one would get from non-perturbative resummations?

# Summary

- Understanding of the mechanisms driving shell evolution (3N forces, tensor forces) allows for a better modelling all over the nuclear chart.
- Pseudo-SU(3) model predictions are in good agreement with SM diagonalizations & SCCM for  $N = 52, 54$  nuclei. Collectivity and triaxiality predicted above  $^{78}\text{Ni}$  core. The same should happen above  $^{132}\text{Sn}$  core.
- No unique multipole interactions for SM from MBPT theory. Could one use on a daily basis TBME interactions from CC or IM-SRG?

Thanks for your attention !