

# Model Stability from Shell far

Frédéric Nowacki<sup>1</sup>



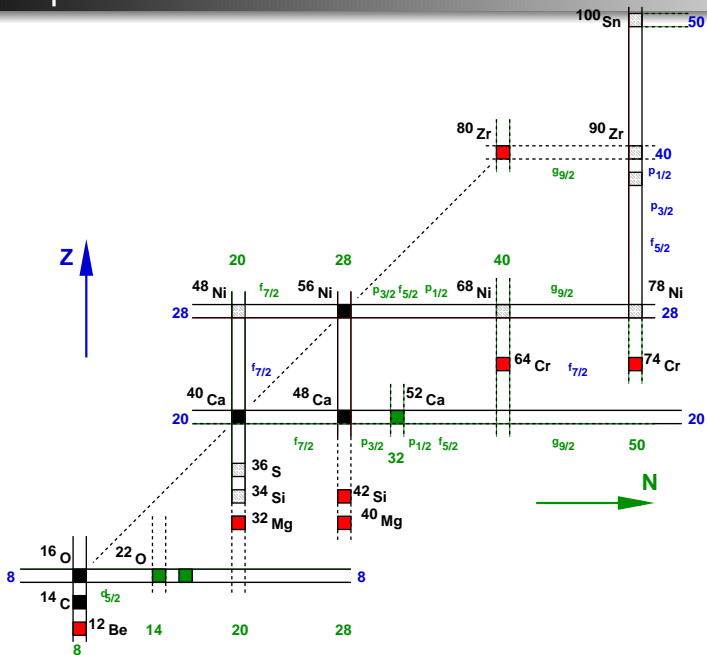
International Collaborations in nuclear theory:  
Theory for open-shell nuclei near the limits of stability

**MICHIGAN STATE**  
**UNIVERSITY**

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<sup>1</sup>Strasbourg-Madrid Shell-Model collaboration

# Landscape of medium mass nuclei



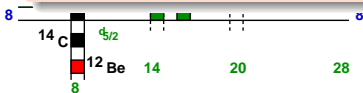
## UNDERSTANDING REGULARITIES for both SPHERICAL and DEFORMED systems

- Magic Numbers:  $^{24}\text{O}$ ,  $^{48}\text{Ni}$ ,  $^{54}\text{Ca}$ ,  $^{78}\text{Ni}$ ,  $^{100}\text{Sn}$
- Islands of Deformation:  $^{12}\text{Be}$ ,  $^{32}\text{Mg}$ ,  $^{42}\text{Si}$ ,  $^{64}\text{Cr}$ ,  $^{80}\text{Zr}$  ...
  
- Variety of phenomena dictated by shell structure
- Close connection between collective behaviour and underlying shell structure
- 

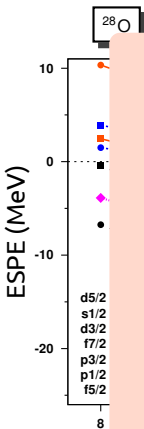
$$\mathcal{H} = \mathcal{H}_m + \mathcal{H}_M$$

Interplay between

- Monopole field (spherical mean field)
- Multipole correlations (pairing, Q.Q, ...)

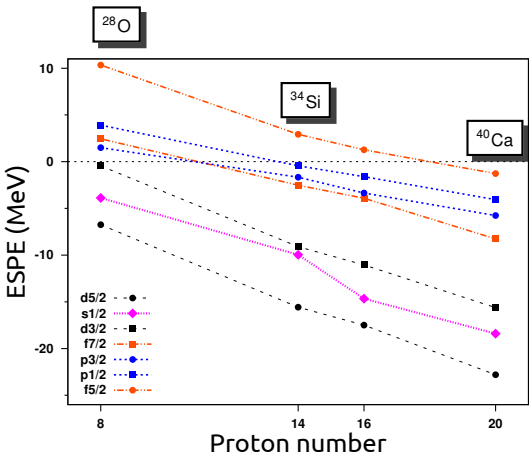


# Effective Single Particle Energies: Trends



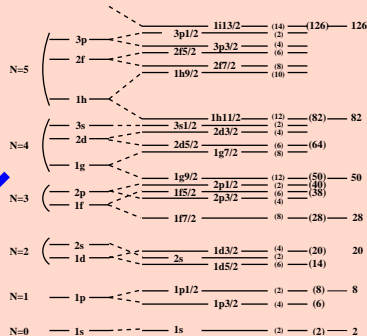
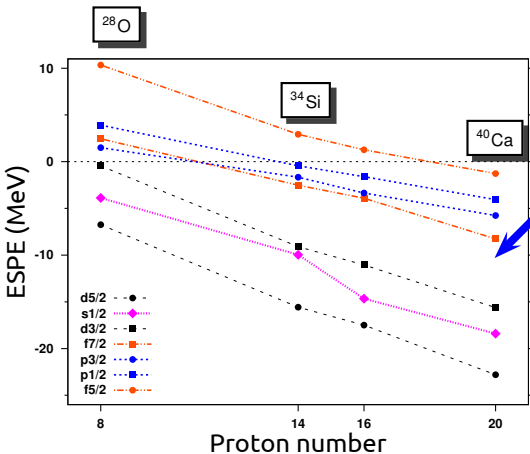
- Magic numbers are associated to energy gaps in the spherical mean field. Therefore, to promote particles above the Fermi levels costs energy
- However some intruders configurations can overwhelm their loss of monopole energy with their huge gain in correlation energy
- Several examples of this phenomenon exist in stable magic nuclei (as in  $^{40}\text{Ca}$  nucleus) in the form of coexisting **spherical**, **deformed** and **superdeformed** states in a very narrow energy range
- At the very neutron rich or very proton rich edges, the  $T=0$  and  $T=1$  channels of the effective nuclear interaction weight very differently than they do at the stability line. Therefore the effective single particle structure may suffer important changes, leading in some cases to the vanishing of established shell closures or to the appearance of new ones

# Effective Single Particle Energies: Trends



- At the neutron drip line, the ESPE's of  $^{28}\text{O}$  are completely at variance with those of  $^{40}\text{Ca}$  at the stability valley. The change from the standard ESPE's of  $^{16}\text{O}$  to the anomalous ones in  $^{28}\text{O}$  is totally due to the interactions of *sd* shell neutrons among themselves
- Notice that the *sd* shell orbits remain always below the *pf* shell with the  $\nu 0f_{7/2}$  and  $\nu 0p_{3/2}$  orbitals DO get inverted
- The monopole part of the neutron-proton interaction restores the N=20 shell gap when the valley of stability is approached
- Spin-Tensor decomposition shows it is mainly a **Central** and **Tensor** effect

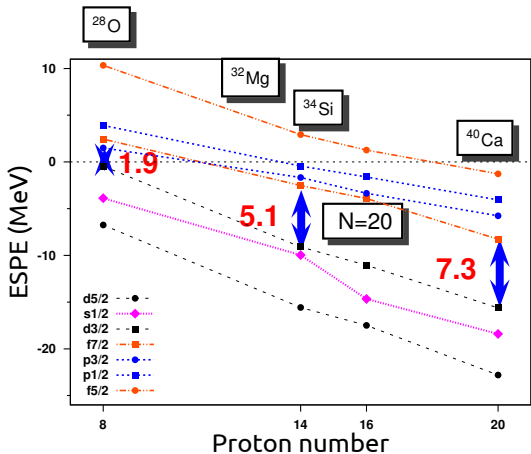
# Effective Single Particle Energies: Trends



N=20 shell gap when the valley of stability is approached

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# Effective Single Particle Energies: Trends



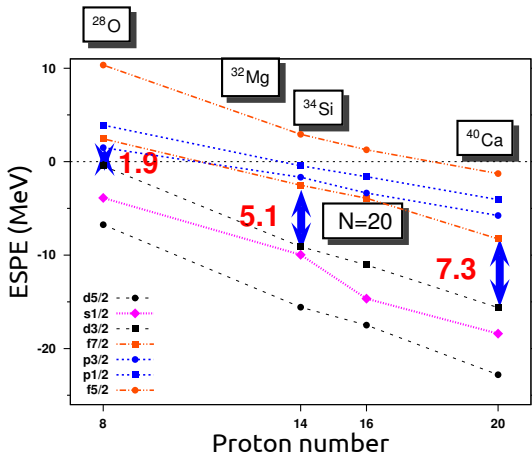
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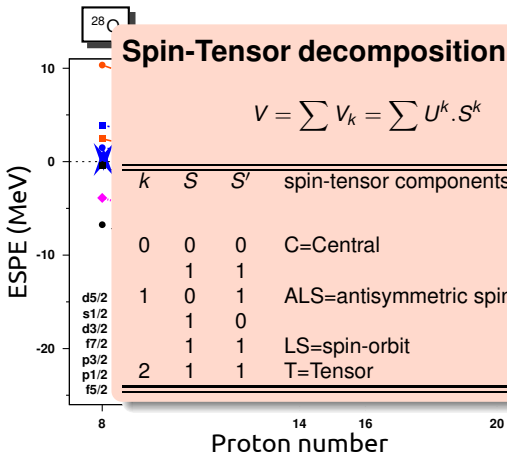
- Notice that the *sd* shell orbits remain always below the *pf* shell with the  $\nu 0f_{7/2}$  and  $\nu 0p_{3/2}$  orbitals DO get inverted

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# Effective Single Particle Energies: Trends



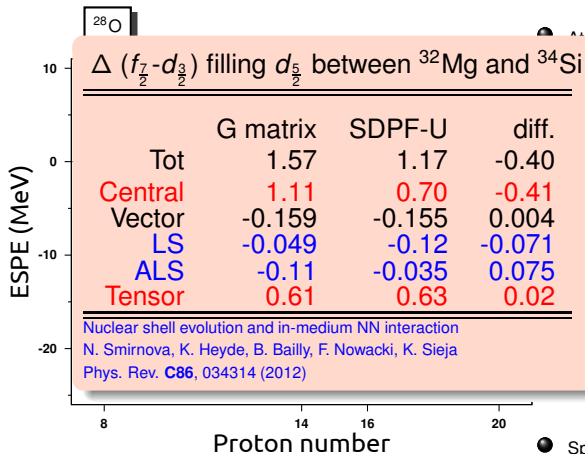
neutron drip line, the ESPE's of  $^{28}\text{Ca}$  are completely at variance with  $^{40}\text{Ca}$  at the stability valley. The anomalous ones in  $^{28}\text{O}$  are due to the interactions of  $sd$  shell among themselves

at the  $sd$  shell orbits remain below the  $pf$  shell with the  $\nu 0f_{7/2}$  orbitals DO get inverted

opole part of the proton interaction restores the shell gap when the valley of  $sd$  is approached

● Spin-Tensor decomposition shows it is mainly a **Central** and **Tensor** effect

# Effective Single Particle Energies: Trends



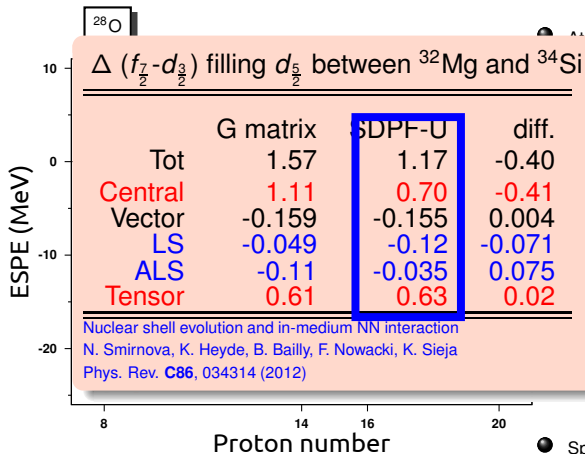
At the neutron drip line, the ESPE's of  $\nu f_{7/2}$  are entirely at variance with the standard ESPE's of  $\nu f_{7/2}$  at the stability valley. The anomalous ones in  $^{28}\text{O}$  is due to the interactions of  $sd$  shell orbitals with themselves

$\nu sd$  shell orbits remain at variance with  $\nu pf$  shell with the  $\nu 0f_{7/2}$  orbitals DO get inverted

A part of the spin-tensor interaction restores the order when the valley of stability is approached

Spin-Tensor decomposition shows it is mainly a **Central** and **Tensor** effect

# Effective Single Particle Energies: Trends



At the neutron drip line, the ESPE's of  $d_{5/2}$  are entirely at variance with the  $d_{3/2}$  at the stability valley. The large negative standard ESPE's of the  $d_{5/2}$  orbital in  $^{28}\text{O}$  is due to the interactions of  $sd$  shell orbitals with themselves

$sd$  shell orbits remain at high energy with  $pf$  shell with the  $\nu 0f_{7/2}$  orbitals DO get inverted

A part of the spin-tensor interaction restores the  $d_{5/2}$  when the valley of  $d_{3/2}$  is approached

Spin-Tensor decomposition shows it is mainly a **Central** and **Tensor** effect

PHYSICAL REVIEW C **79**, 034318 (2009)

## Nuclear “bubble” structure in $^{34}\text{Si}$

M. Grasso,<sup>1</sup> L. Gaudefroy,<sup>2</sup> E. Khan,<sup>1</sup> T. Nikšić,<sup>3</sup> J. Piekarewicz,<sup>4</sup> O. Sorlin,<sup>5</sup> N. Van Giai,<sup>1</sup> and D. Vretenar<sup>3</sup>

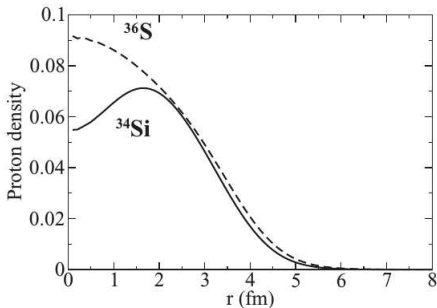
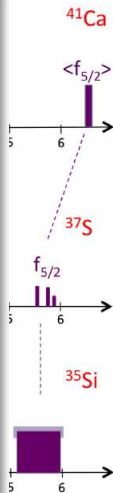
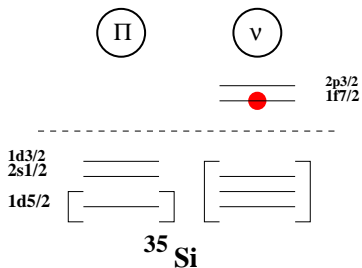


FIG. 4. HF charge densities (in units of  $\text{fm}^{-3}$ ) of  $^{36}\text{S}$  (dashed line) and  $^{34}\text{Si}$  (solid line) calculated with the Skyrme interaction SLy4.

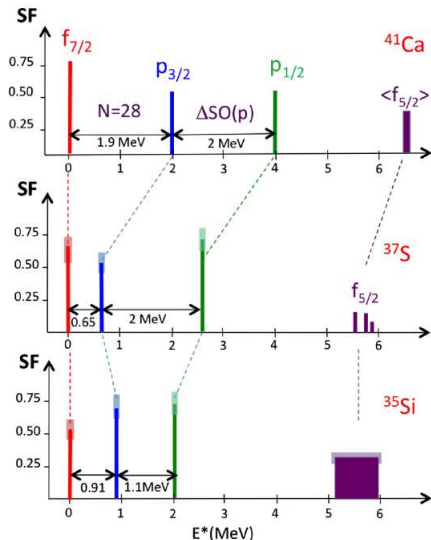


# Effective Single Particle Energies: Trends



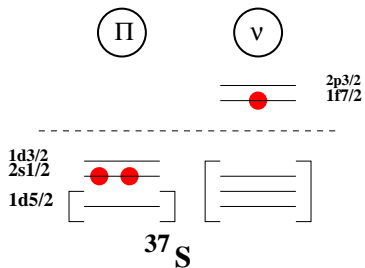
$\Delta (p_{3/2} - p_{1/2})$  filling  $s_{1/2}$  between  $^{34}\text{Si}$  and  $^{36}\text{S}$

	KLS	N3LO (bare)	N3LO ( $4\hbar\omega$ )
Tot	0.58	0.58	0.60
Central	0.00	0.00	0.00
<b>Vector</b>	<b>0.58</b>	<b>0.58</b>	<b>0.60</b>
Tensor	0.00	0.00	0.00



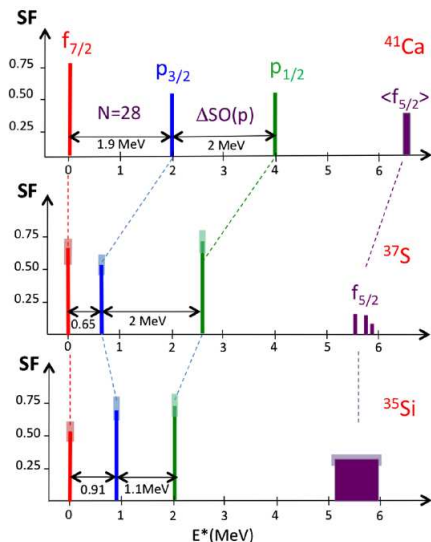
G. Burgunder, PhD Thesis  
GANIL, December 2011

# Effective Single Particle Energies: Trends



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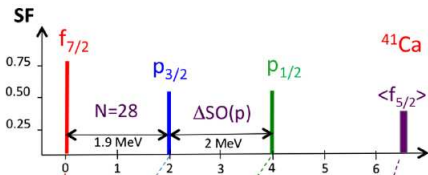


G. Burgunder, PhD Thesis  
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# Effective Single Particle Energies: Trends



2p<sub>3/2</sub>  
1i<sub>7/2</sub>



PRL 112, 042502 (2014)

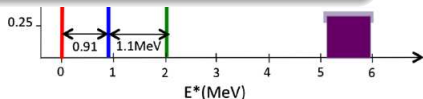
PHYSICAL REVIEW LETTERS

week ending  
31 JANUARY 2014

## Experimental Study of the Two-Body Spin-Orbit Force in Nuclei

G. Burgunder,<sup>1</sup> O. Sorlin,<sup>1</sup> F. Nowacki,<sup>2</sup> S. Giron,<sup>3</sup> F. Hammache,<sup>3</sup> M. Moukaddam,<sup>2</sup> N. de Séréville,<sup>3</sup> D. Beaulieu,<sup>3</sup> L. Càceres,<sup>1</sup> E. Clément,<sup>1</sup> G. Duchêne,<sup>2</sup> J. P. Ebran,<sup>4</sup> B. Fernandez-Dominguez,<sup>1,5</sup> F. Flavigny,<sup>6</sup> S. Franchou,<sup>3</sup> J. Gibelin,<sup>7</sup> A. Gillibert,<sup>6</sup> S. Grévy,<sup>1,8</sup> J. Guillot,<sup>3</sup> A. Lepailleur,<sup>1</sup> I. Matea,<sup>3</sup> A. Matta,<sup>3,9</sup> L. Nalpas,<sup>6</sup> A. Obertelli,<sup>6</sup> T. Otsuka,<sup>10</sup> J. Pancin,<sup>1</sup> A. Poves,<sup>11</sup> R. Raabe,<sup>1,12</sup> J. A. Scarpaci,<sup>13</sup> I. Stefan,<sup>3</sup> C. Stodel,<sup>1</sup> T. Suzuki,<sup>14</sup> and J. C. Thomas<sup>1</sup>

	( <sup>41</sup> Ca)	( <sup>41</sup> S)	( <sup>41</sup> Si)
Tot	0.58	0.58	0.60
Central	0.00	0.00	0.00
Vector	0.58	0.58	0.60
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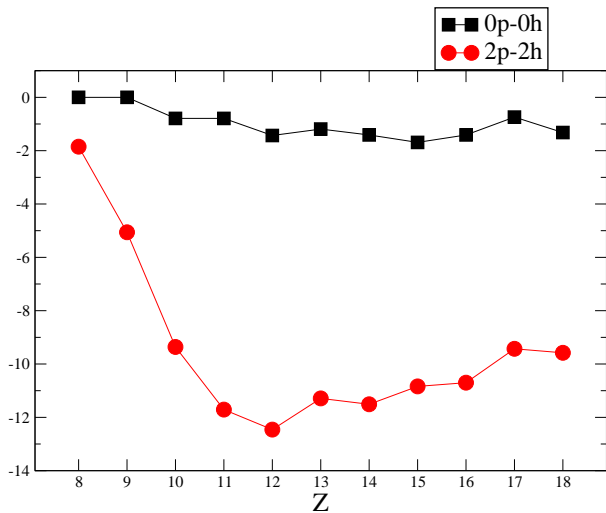
G. Burgunder, PhD Thesis  
GANIL, December 2011

# Correlation Energies

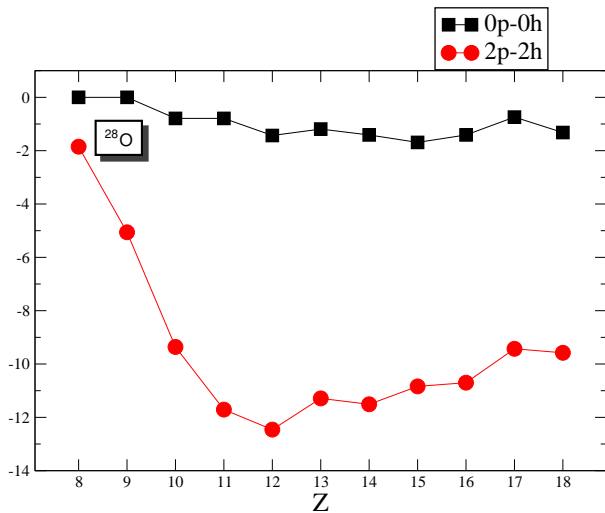
- Let us consider the configurations with closed  $N=20$   $[sd]^{12,Z}$  (normal filling) and the ones with two neutrons blocked in the  $pf$  shell  $[sd]^{10,Z}[pf]^{2,0}$  (intruder)
- And calculate the energy of the ground states at fixed configuration, with and without correlations



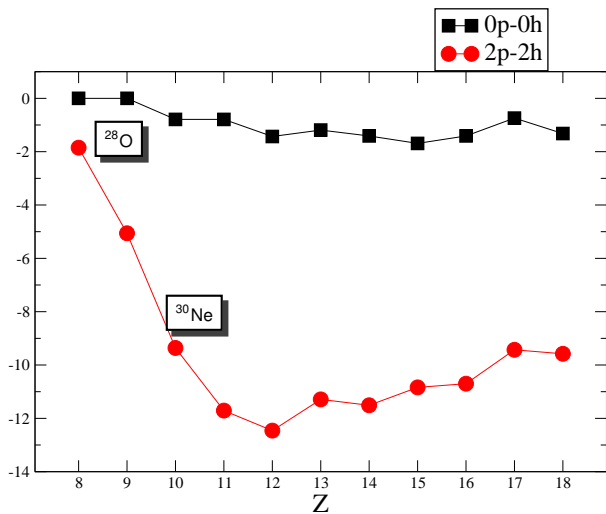
# Correlations energies: normal vs 2p-2h intruder



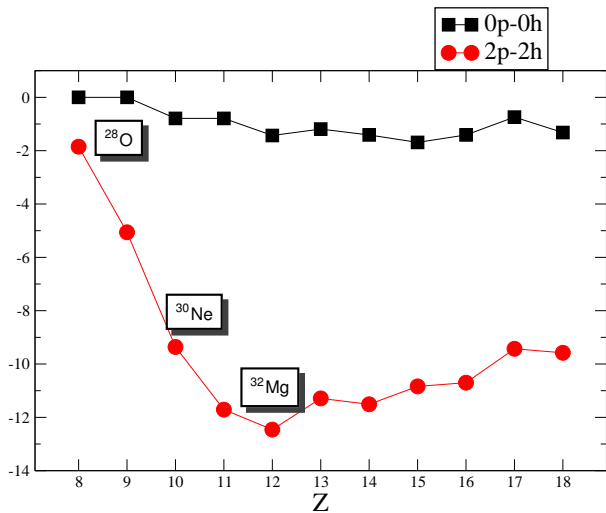
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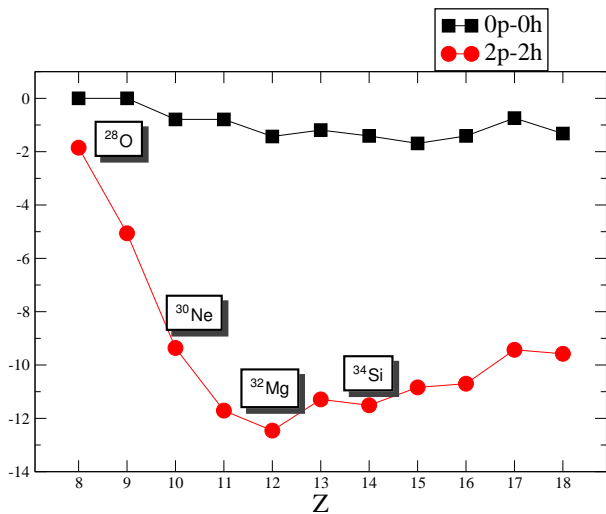
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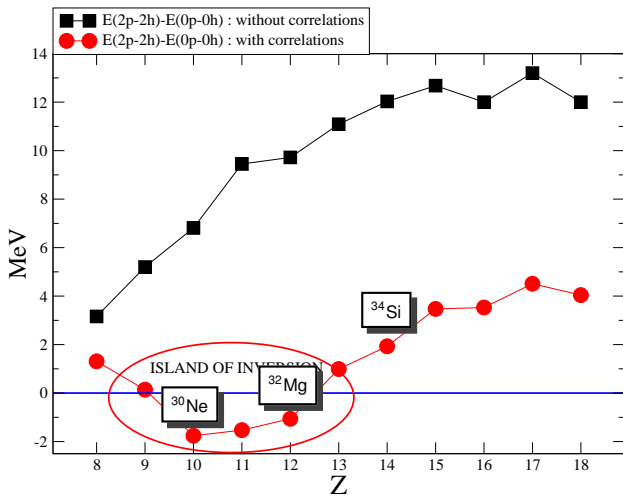
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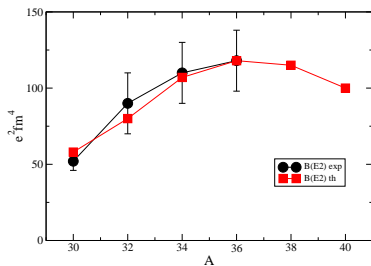
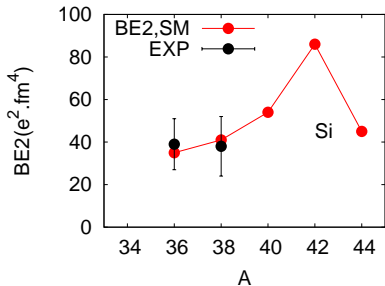
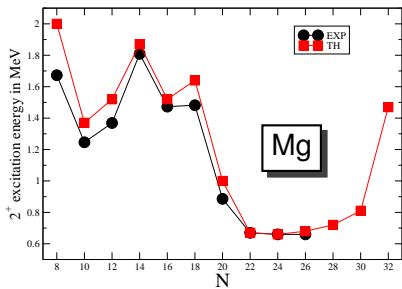
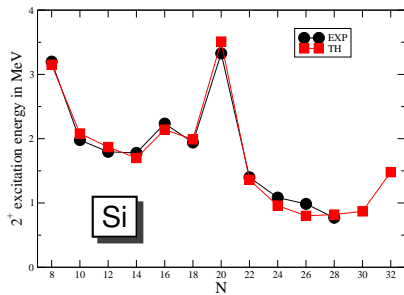
# Correlations energies: normal vs 2p-2h intruder



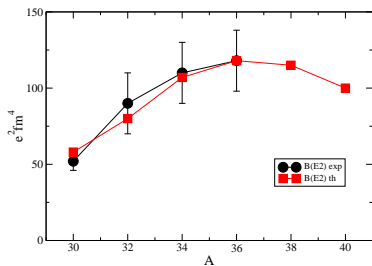
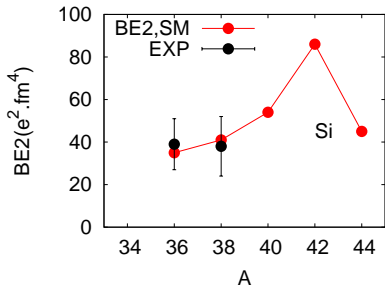
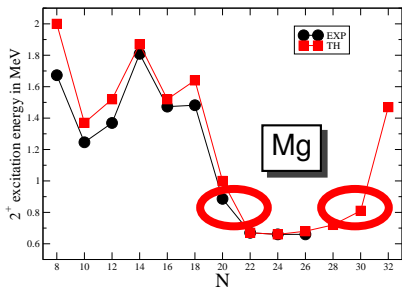
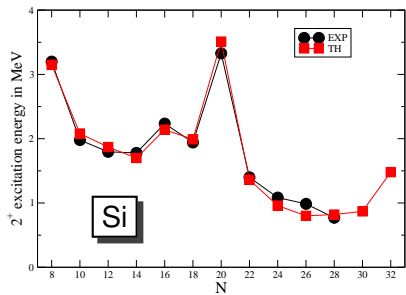
# Gaps: normal vs 2p-2h intruder



# Merging of the islands of inversion at $N=20$ and $N=28$

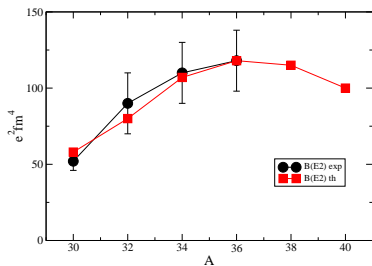
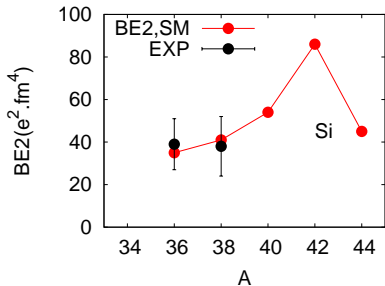
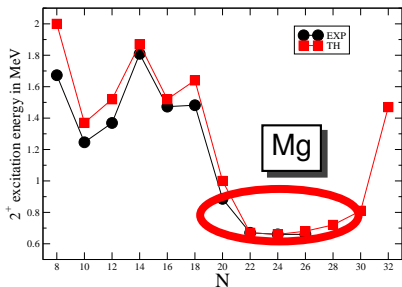
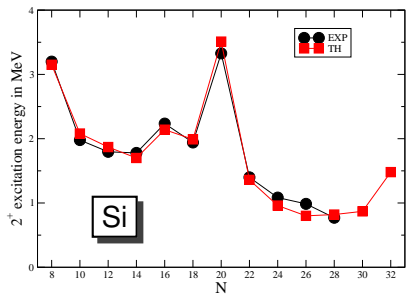


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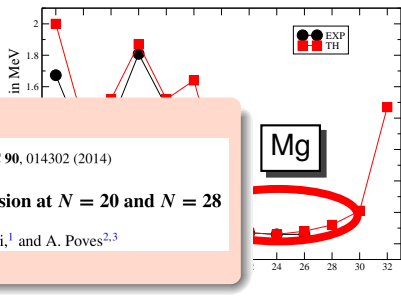
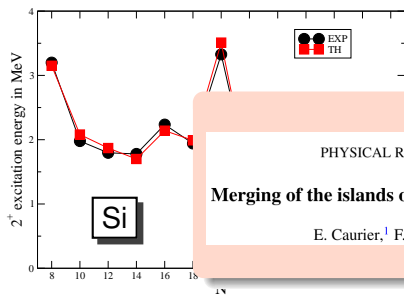




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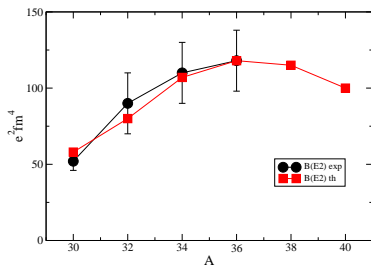
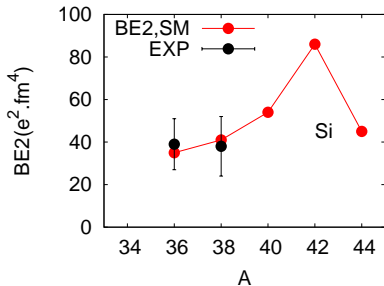
# Merging of the islands of inversion at $N=20$ and $N=28$



PHYSICAL REVIEW C **90**, 014302 (2014)

**Merging of the islands of inversion at  $N = 20$  and  $N = 28$**

E. Caurier,<sup>1</sup> F. Nowacki,<sup>1</sup> and A. Poves<sup>2,3</sup>



ER and FE AROUND  $N=40$

A NEW REGION OF DEFORMATION.

## A. Poves

$E4_{G}$

$$g(0p4 - 2p4) = 5.70$$

$$g(0p4 - 4p4) = 8.30$$

$$Q = -9.0 b^2$$

$$CS < 1\%$$

$$BE2 = 19.8 b^4$$

$$u(d5/2) = 1.1$$

$$\frac{E(4^+)}{E(2^+)} = 2.7$$

$$\left[ \frac{E(4^+)}{E(2^+)} = (3.2)(3.4) \right]$$

in the intruder configurations.

A SITUATION THAT REMINDS WHAT IS KNOWN AT  $N=20$  FFS.

# Island of inversion at N=40, an old story: 2003

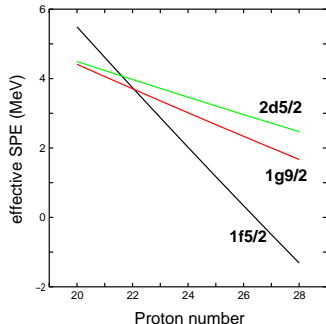
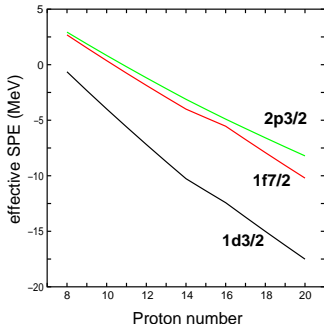
Eur. Phys. J. **16**, 55–61 (2003)  
DOI 10.1140/epj/i2002-10069-9

THE EUROPEAN  
PHYSICAL JOURNAL A

GANIL

## New region of deformation in the neutron-rich ${}^{60}_{24}\text{Cr}_{36}$ and ${}^{62}_{24}\text{Cr}_{38}$

O. Sorlin<sup>1,a</sup>, C. Donzaud<sup>1</sup>, F. Nowacki<sup>2</sup>, J.C. Angélique<sup>3</sup>, F. Azaiez<sup>1</sup>, C. Bourgeois<sup>1</sup>, V. Chisté<sup>1</sup>, Z. Dlouhy<sup>4</sup>, S. Grévy<sup>3</sup>, D. Guillemaud-Mueller<sup>1</sup>, F. Ibrahim<sup>1</sup>, K.-L. Kratz<sup>5</sup>, M. Lewitowicz<sup>6</sup>, S.M. Lukyanov<sup>7</sup>, J. Mrasek<sup>4</sup>, Yu.-E. Penionzhkevich<sup>7</sup>, F. de Oliveira Santos<sup>6</sup>, B. Pfeiffer<sup>5</sup>, F. Pougheon<sup>1</sup>, A. Poves<sup>8</sup>, M.G. Saint-Laurent<sup>6</sup>, and M. Stanoiu<sup>6</sup>



MSU

RAPID COMMUNICATIONS

PHYSICAL REVIEW C **81**, 051304(R) (2010)

## Collectivity at $N = 40$ in neutron-rich $^{64}\text{Cr}$

A. Gade,<sup>1,2</sup> R. V. F. Janssens,<sup>3</sup> T. Baugher,<sup>1,2</sup> D. Bazin,<sup>1</sup> B. A. Brown,<sup>1,2</sup> M. P. Carpenter,<sup>3</sup> C. J. Chiara,<sup>3,4</sup> A. N. Deacon,<sup>5</sup> S. J. Freeman,<sup>5</sup> G. F. Grinyer,<sup>1</sup> C. R. Hoffman,<sup>3</sup> B. P. Kay,<sup>3</sup> F. G. Kondev,<sup>6</sup> T. Lauritsen,<sup>3</sup> S. McDaniel,<sup>1,2</sup> K. Meierbachtol,<sup>1,7</sup> A. Ratkiewicz,<sup>1,2</sup> S. R. Stroberg,<sup>1,2</sup> K. A. Walsh,<sup>1,2</sup> D. Weisshaar,<sup>1</sup> R. Winkler,<sup>1</sup> and S. Zhu<sup>3</sup>

<sup>1</sup>National Superconducting Cyclotron Laboratory, Michigan State University, East Lansing, Michigan 48824, USA

<sup>2</sup>Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824, USA

<sup>3</sup>Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

<sup>4</sup>Department of Physics, University of Illinois at Chicago, Chicago, Illinois 60607, USA

GANIL

RAPID COMMUNICATION

PHYSICAL REVIEW C **81**, 061301(R) (2010)

## Onset of collectivity in neutron-rich Fe isotopes: Toward a new island of inversion?

J. Ljungvall,<sup>1,2,3</sup> A. Görgen,<sup>1</sup> A. Obertelli,<sup>1</sup> W. Korten,<sup>1</sup> E. Clément,<sup>2</sup> G. de France,<sup>2</sup> A. Bürger,<sup>4</sup> J.-P. Delaroche,<sup>5</sup> A. Dewald,<sup>6</sup> A. Gadea,<sup>7</sup> L. Gaudefroy,<sup>5</sup> M. Girod,<sup>5</sup> M. Hackstein,<sup>6</sup> J. Libert,<sup>8</sup> D. Mengoni,<sup>9</sup> F. Nowacki,<sup>10</sup> T. Pissulla,<sup>6</sup> A. Poves,<sup>11</sup> F. Recchia,<sup>12</sup> M. Rejmund,<sup>2</sup> W. Rother,<sup>6</sup> E. Sahin,<sup>12</sup> C. Schmitt,<sup>2</sup> A. Shrivastava,<sup>2</sup> K. Sieja,<sup>10</sup> J. J. Valiente-Dobón,<sup>12</sup> K. O. Zell,<sup>6</sup> and M. Zielińska<sup>13</sup>

<sup>1</sup>CEA Saclay, IRFU, Service de Physique Nucléaire, F-91191 Gif-sur-Yvette, France

<sup>2</sup>GANIL, CEA/DSM-CNRS/IN2P3, Bd Henri Becquerel, BP 55027, F-14076 Caen, France

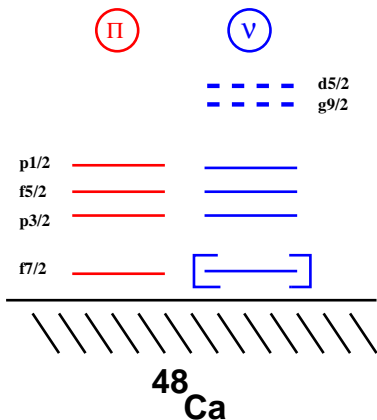
<sup>3</sup>CEA/DSM-CNRS/IN2P3, F-91191 Gif-sur-Yvette, France



Island of inversion around  $^{64}\text{Cr}$

S. Lenzi, F. Nowacki, A. Poves and K. Sieja

Phys. Rev. C 82, 054301, 2010.



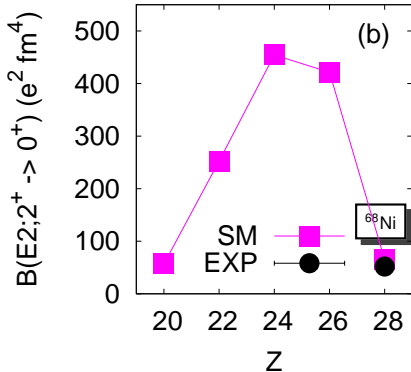
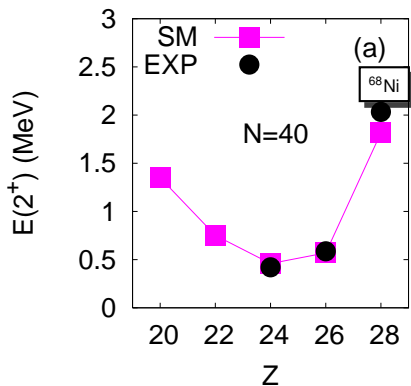
## LNPS interaction:

- based on realistic TBME
- new fit of the pf shell (KB3GR, E. Caurier)
- monopole corrections
- $g_{9/2}$ - $d_{5/2}$  gap set to 1.5 Mev in  $^{68}\text{Ni}$

## Calculations:

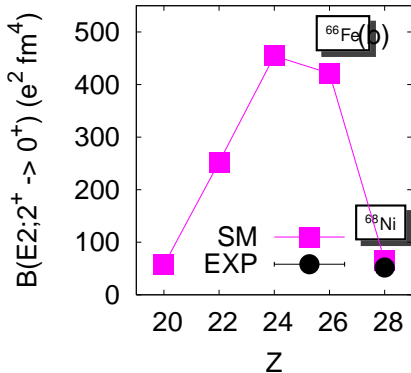
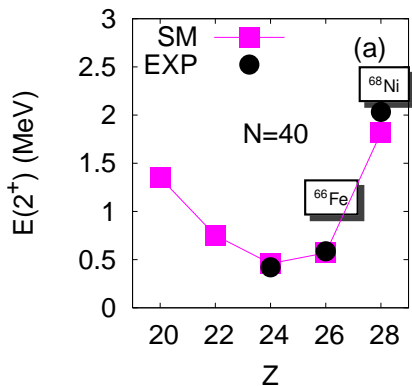
- up to 14p-14h excitations across  $Z=28$  and  $N=40$  gaps
- up to  $10^{10}$
- m-scheme code ANTOINE (non public version)

# Shape transition at N=40



Nucleus	$\nu g_{9/2}$	$\nu d_{5/2}$	configuration
$^{68}\text{Ni}$	0.98	0.10	0p0h(51%)
$^{66}\text{Fe}$	3.17	0.46	4p4h(26%)
$^{64}\text{Cr}$	3.41	0.76	6p6h(23%)
$^{62}\text{Ti}$	3.17	1.09	4p4h(48%)

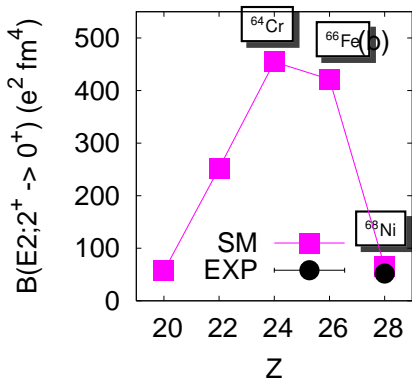
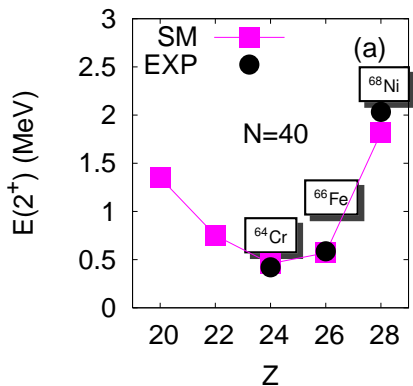
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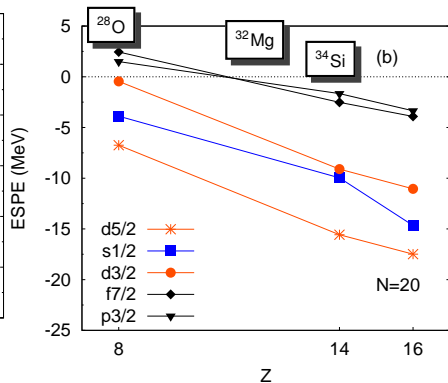
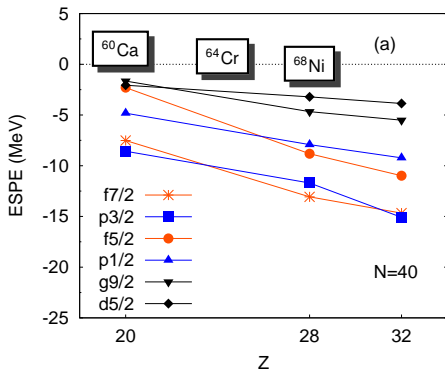


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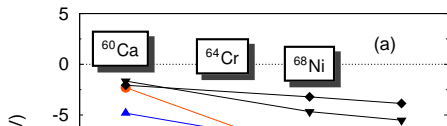
# Neutron effective single particle energies



- reduction of the  $\nu f_{5/2}$ - $g_{9/2}$  gap with removing  $f_{7/2}$  protons
- proximity of the quasi-SU3 partner  $d_{5/2}$
- inversion of  $d_{5/2}$  and  $g_{9/2}$  orbitals  
same ordering as CC calculations

- reduction of the  $\nu d_{3/2}$ - $f_{7/2}$  gap with removing  $d_{5/2}$  protons
- proximity of the quasi-SU3 partner  $p_{3/2}$
- inversion of  $p_{3/2}$  and  $f_{7/2}$  orbitals

# Neutron effective single particle energies



PRL **109**, 032502 (2012)

PHYSICAL REVIEW LETTERS

TABLE II. Energies of the  $5/2^+$  and  $9/2^+$  resonances in  $^{53,55,61}\text{Ca}$ .  $\text{Re}[E]$  is the energy relative to the one-neutron emission threshold, and the width is  $\Gamma = -2\text{Im}[E]$  (in MeV).

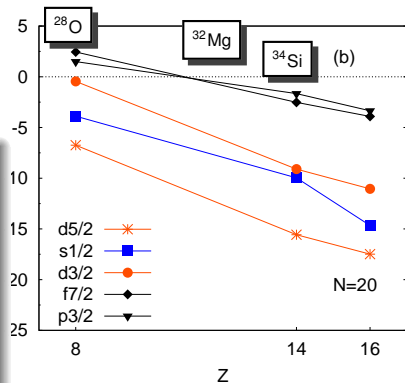
	$^{53}\text{Ca}$		$^{55}\text{Ca}$		$^{61}\text{Ca}$	
$J^\pi$	$\text{Re}[E]$	$\Gamma$	$\text{Re}[E]$	$\Gamma$	$\text{Re}[E]$	$\Gamma$
$5/2^+$	1.99	1.97	1.63	1.33	1.14	0.62
$9/2^+$	4.75	0.28	4.43	0.23	2.19	0.02

G. Hagen et al.

Phys. Rev. Lett. **109**, 032502 (2012)

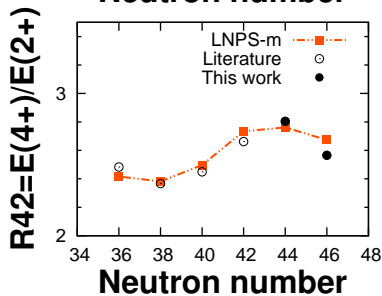
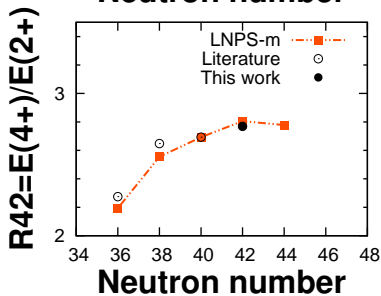
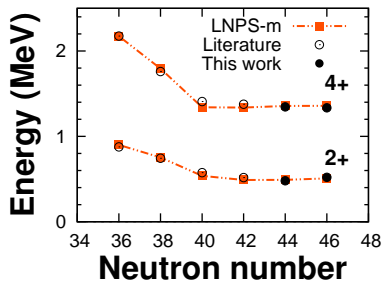
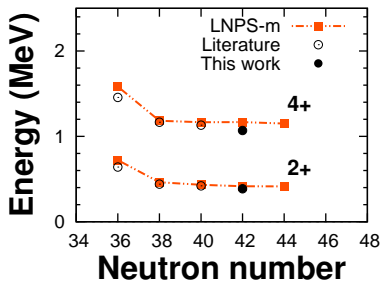
removing  $f_{7/2}$  protons

- proximity of the quasi-SU3 partner  $d_{5/2}$
- inversion of  $d_{5/2}$  and  $g_{9/2}$  orbitals  
same ordering as CC calculations



- reduction of the  $\nu d_{3/2}-f_{7/2}$  gap with removing  $d_{5/2}$  protons
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# Extension of collectivity N=40 towards N=50



# Extension of collectivity N=40 towards N=50

Energy (MeV)

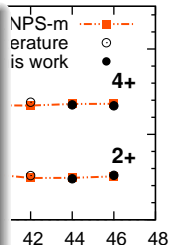
TABLE I: Quadrupole deformation properties of Cr and Fe isotopes. Energies are in MeV,  $B(E2)$  in  $e^2 \text{ fm}^4$ , and  $Q$  in  $e \text{ fm}^2$ . Experimental energies are the same as Fig. 3.

	$^{62}\text{Cr}$	$^{64}\text{Cr}$	$^{66}\text{Cr}$	$^{68}\text{Cr}$	$^{66}\text{Fe}$	$^{68}\text{Fe}$	$^{70}\text{Fe}$	$^{72}\text{Fe}$
$E^*(2_1^+)$ exp.	0.44	0.42	0.39	-	0.57	0.52	0.48	0.52
$E^*(2_1^+)$ theo.	0.46	0.43	0.42	0.41	0.54	0.49	0.49	0.51
$Q_{spec}$	-38	-38	-39	-38	-37	-40	-39	-33
$B(E2)\downarrow$ th.	378	388	389	367	372	400	382	279
$Q_{int}$ from $Q_{spec}$	135	136	137	132	131	140	135	116
$Q_{int}$ from $B(E2)$	138	140	140	136	137	142	139	118
$\langle \beta \rangle$	0.33	0.33	0.32	0.30	0.29	0.30	0.28	0.24
$E^*(4_1^+)$ exp.	1.17	1.13	1.07	-	1.41	1.39	1.35	1.33
$E^*(4_1^+)$ theo.	1.18	1.13	1.06	1.15	1.34	1.34	1.36	1.36
$Q_{spec}$	-49	-49	-46	-47	-47	-51	-48	-40
$B(E2)\downarrow$ th.	562	534	562	530	553	608	574	377
$Q_{int}$ from $Q_{spec}$	135	134	134	130	129	141	132	111
$Q_{int}$ from $B(E2)$	141	140	141	137	139	146	142	115
$\langle \beta \rangle$	0.34	0.33	0.32	0.31	0.29	0.30	0.29	0.23

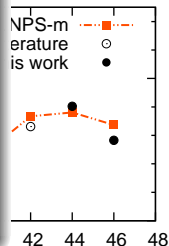
$R42=E(4+)/E(2+)$

Neutron number

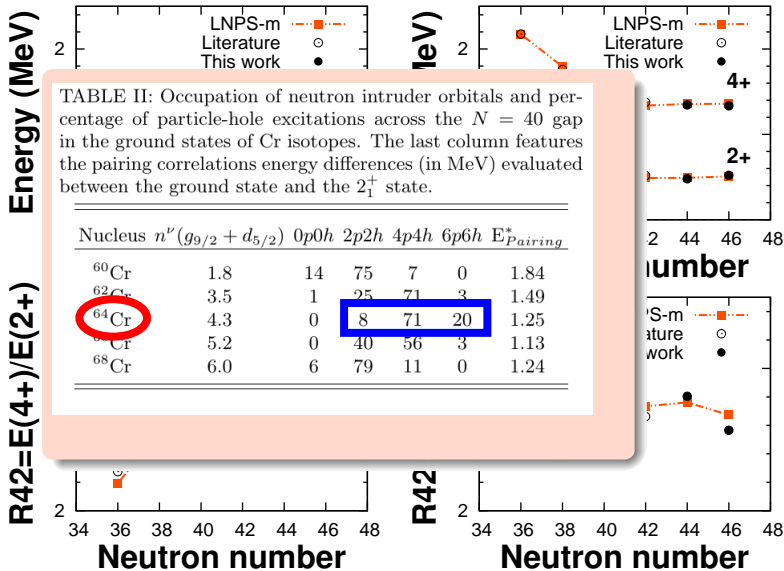
Neutron number



number



# Extension of collectivity N=40 towards N=50

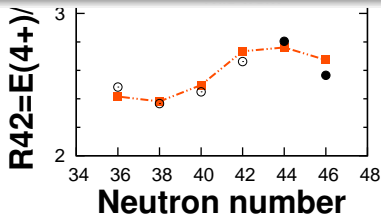
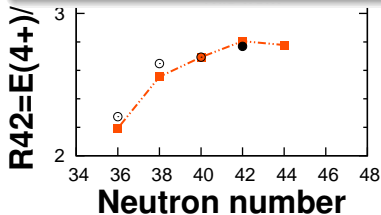


# Extension of collectivity N=40 towards N=50

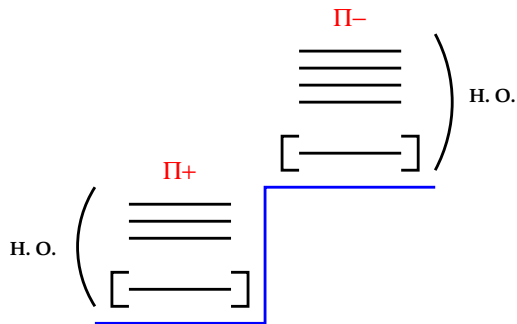
## Extension of the N=40 Island of Inversion towards N=50: Spectroscopy of $^{66}\text{Cr}$ , $^{70,72}\text{Fe}$

C. Santamaria,<sup>1,2</sup> C. Louchart,<sup>3</sup> A. Obertelli,<sup>1,2</sup> V. Werner,<sup>3,4</sup> P. Doornenbal,<sup>2</sup> F. Nowacki,<sup>5</sup> G. Authelet,<sup>1</sup> H. Baba,<sup>2</sup> D. Calvet,<sup>1</sup> F. Château,<sup>1</sup> A. Corsi,<sup>1</sup> A. Delbart,<sup>1</sup> J.-M. Gheller,<sup>1</sup> A. Gillibert,<sup>1</sup> T. Isobe,<sup>2</sup> V. Lapoux,<sup>1</sup> M. Matsushita,<sup>6</sup> S. Momiyama,<sup>2,7</sup> T. Motobayashi,<sup>2</sup> M. Niikura,<sup>7</sup> H. Otsu,<sup>2</sup> C. Péron,<sup>1</sup> A. Peyaud,<sup>1</sup> F.C. Pollacco,<sup>1</sup> J.-Y. Roussé,<sup>1</sup> H. Sakurai,<sup>2,7</sup> M. Sasano,<sup>2</sup> Y. Shiga,<sup>2,8</sup> S. Takeuchi,<sup>2</sup> R. Taniuchi,<sup>2,7</sup> T. Taniuchi,<sup>2</sup> H. Wang,<sup>2</sup> K. Yoneda,<sup>2</sup> F. Browne,<sup>9</sup> L.X. Chung,<sup>10</sup> Zs. Dombradi,<sup>11</sup> S. Franchino,<sup>11</sup> M. Giacoppo,<sup>13</sup> A. Gottardo,<sup>12</sup> K. Hadynska-Klek,<sup>13</sup> Z. Korkulu,<sup>11</sup> S. Koyama,<sup>2,7</sup> Y. Koyama,<sup>11</sup> W. Lee,<sup>14</sup> M. Lettmann,<sup>3</sup> R. Lozeva,<sup>5</sup> K. Matsui,<sup>2,7</sup> T. Miyazaki,<sup>2,7</sup> S. Nishimura,<sup>2</sup> J. Nishinaka,<sup>11</sup> S. Ota,<sup>6</sup> Z. Patel,<sup>15</sup> N. Pietralla,<sup>3</sup> E. Sahin,<sup>13</sup> C. Shand,<sup>15</sup> P.-A. Söderström,<sup>2</sup> J. Stefan,<sup>12</sup> C. Steppenbeck,<sup>6</sup> T. Sumikama,<sup>16</sup> D. Suzuki,<sup>12</sup> Zs. Vajta,<sup>11</sup> J. Wu,<sup>2,17</sup> and Z. Xu<sup>14</sup>

c. santamaria et al., submitted for publication in Phys. Rev. Lett.



# Spin-orbit shell closure far from stability



- sd-pf:  $^{42}\text{Si}$  deformed

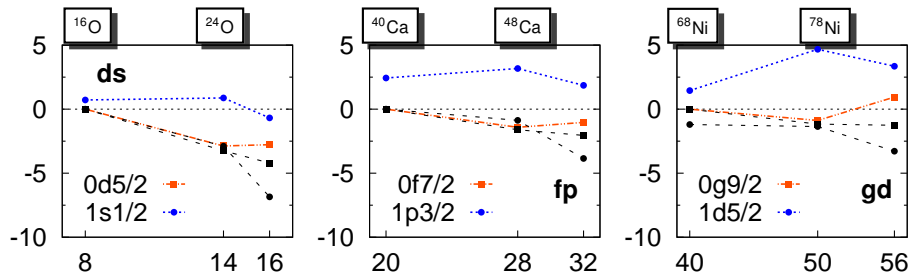
- pf-sdg:  $^{78}\text{Ni}$  ???

- sdg-phf:  $^{132}\text{Sn}$   
doubly magic

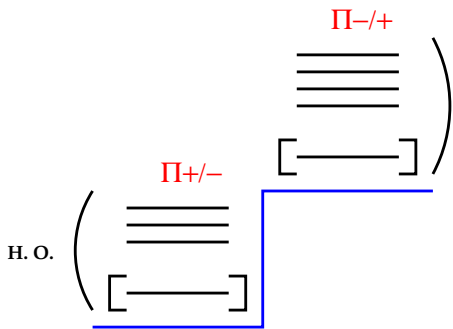
- Evolution of  $Z=28$  from  $N=40$  to  $N=50$
- Evolution of  $N=50$  from  $Z=40$  to  $Z=28$



# Three-body forces in medium mass nuclei



- Evolution of the neutron effective single-particle energies with neutron filling in ds, fp, and gd shells
- “Universal” mechanism for the generation of T=1 spin-orbit shell closures
- Connection with 3N forces and ab-initio calculations:
  - “works” for Coupled-Cluster and to be checked in nickels
  - problems for “ab-initio” core shell-model

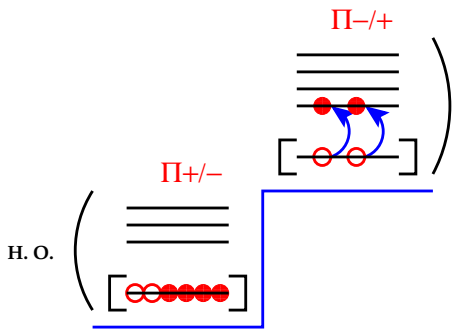


## PFSDG-U interaction:

- based on realistic TBME
- pf shell for protons and gds shell for neutrons
- monopole corrections
- H. O. proton and neutrons gap  $^{78}\text{Ni}$  fixed to phenomenological derived values

## Calculations:

- excitations across  $Z=28$  and  $N=50$  gaps
- up to  $10^{10}$  Slater Determinant basis states
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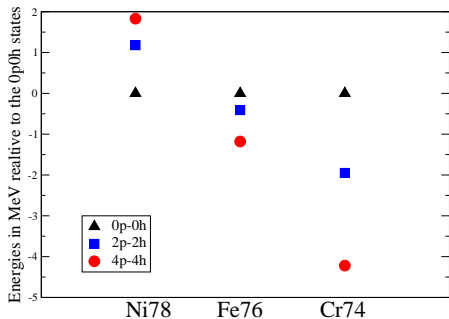
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## Schematic SU3 predictions: [arXiv 1404.0224](https://arxiv.org/abs/1404.0224)

Nilsson-SU3 selfconsistency:  
Quadrupole dominance in heavy N=Z nuclei



- monopole + quadrupole model
- proton gap (5MeV) and neutron gap (5 MeV) estimates
- Quasi-SU3 (protons) and Pseudo-SU3 (neutrons) blocks

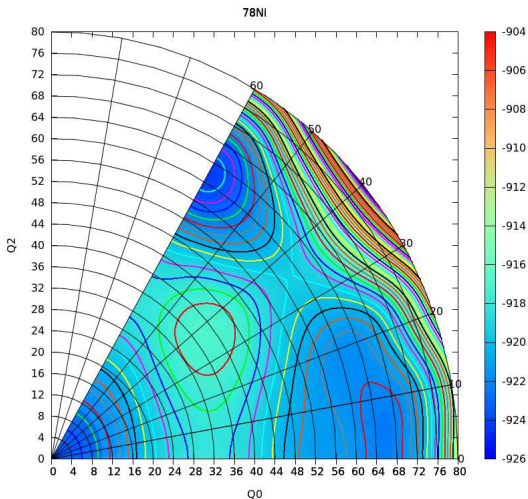
$$Q_s = (\langle 2q_{20} \rangle + 3.)b^2/3.5$$

$$E = \epsilon_i \langle n_i \rangle - \hbar\omega\kappa \left( \frac{\langle 2q_{20(3)} \rangle}{15} + \frac{\langle 2q_{20(4)} \rangle}{23} \right)$$

- Prediction of Island of strong collectivity below  $^{78}\text{Ni}$  !!!

# Shape coexistence in $^{78}\text{Ni}$

- At first approximation,  $^{78}\text{Ni}$  has a double closed shell structure for GS
- But very low-lying  $4p4h$  structures
- The first excited state  $2_1^+$  in  $^{78}\text{Ni}$  predicted at 2.8 MeV and to be a deformed intruder !!!
- Necessity to go beyond  $(fp_{\frac{9}{2}} d_{\frac{5}{2}})$  LNPS space

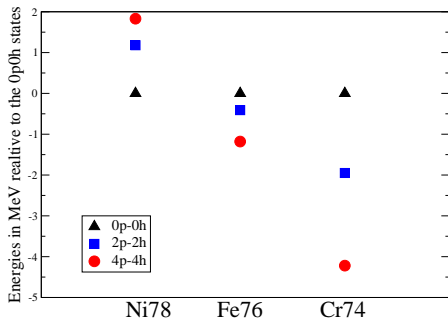


Constrained deformed HF in the SM basis

(B. Bounthong, Ph D Thesis, Strasbourg)

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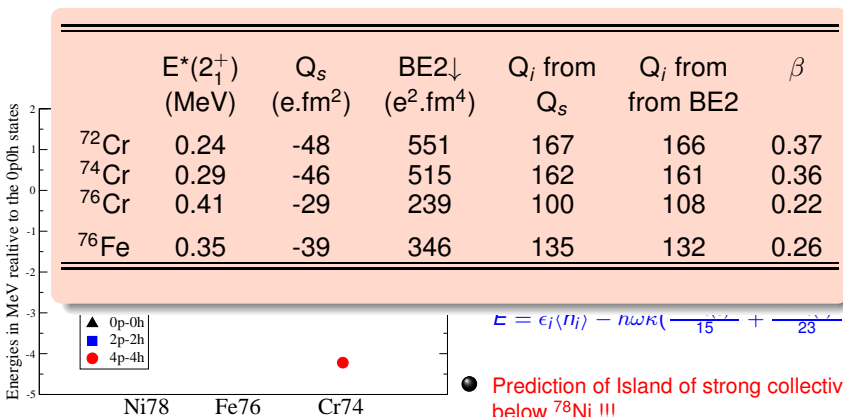
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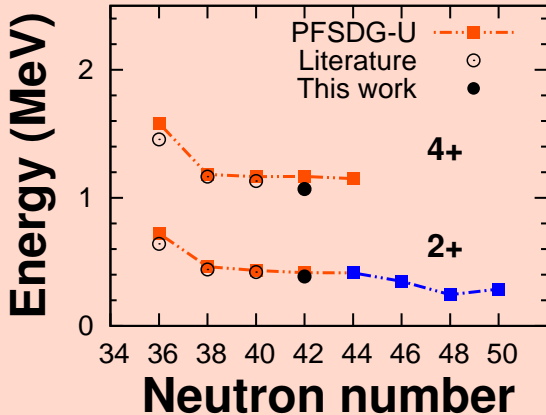
Schematic SU3 predictions: [arXiv 1404.0224](https://arxiv.org/abs/1404.0224)

Nilsson-SU3 selfconsistency:  
Quadrupole dominance in heavy N=Z nuclei



# Island of Deformation below $^{78}\text{Ni}$

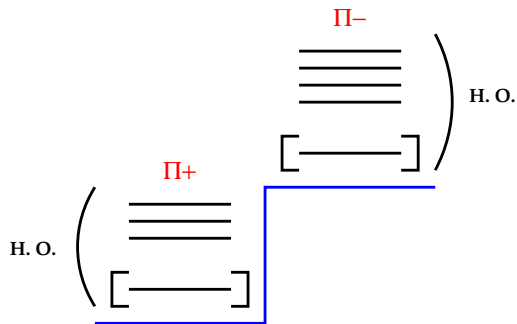
Schematic SU3 predictions: [arXiv 1404.0224](https://arxiv.org/abs/1404.0224)



Extension of Island of deformation to N=50 !!!



# Spin-orbit shell closure far from stability



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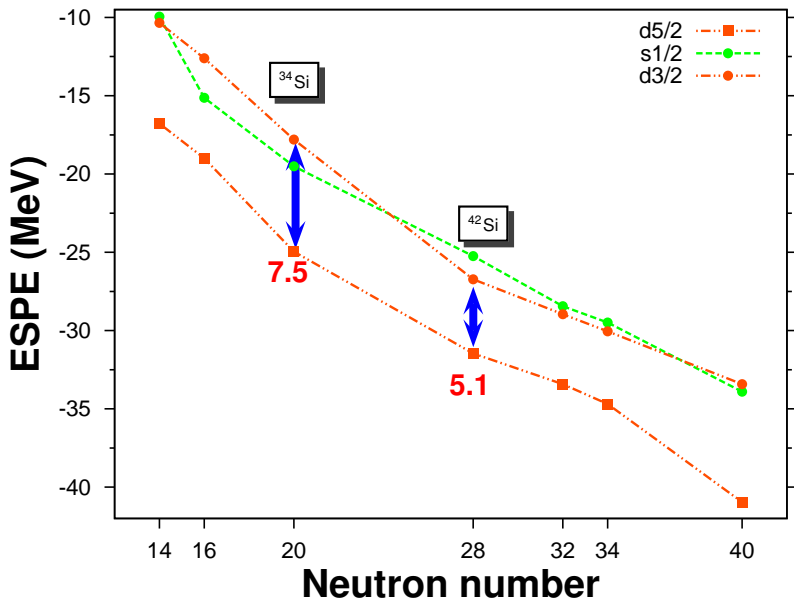
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- Evolution of  $N=50$  from  $Z=40$  to  $Z=28$

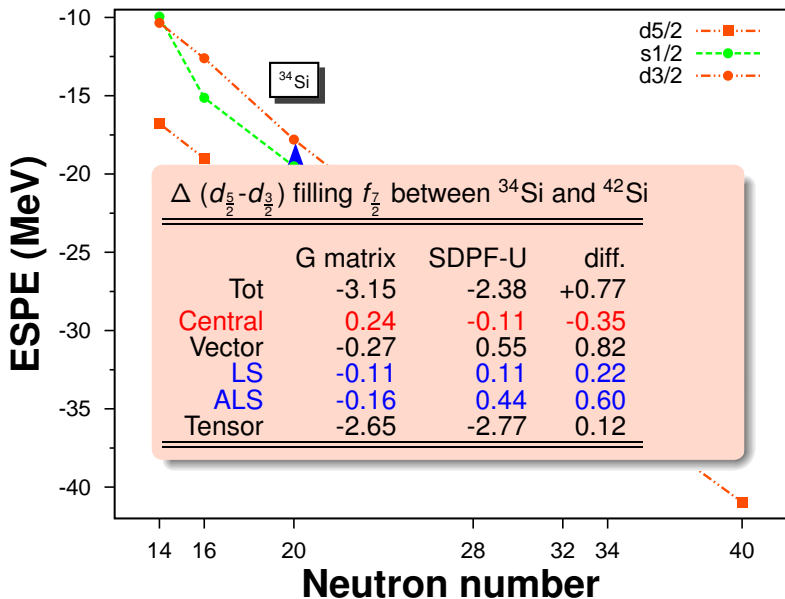
# Effective Single Particle Energies: Trends

## Silicium chain

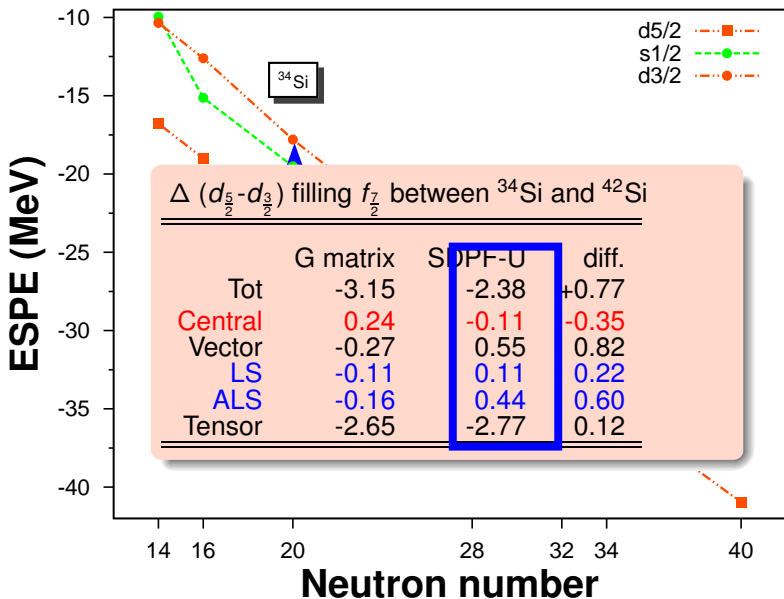


# Effective Single Particle Energies: Trends

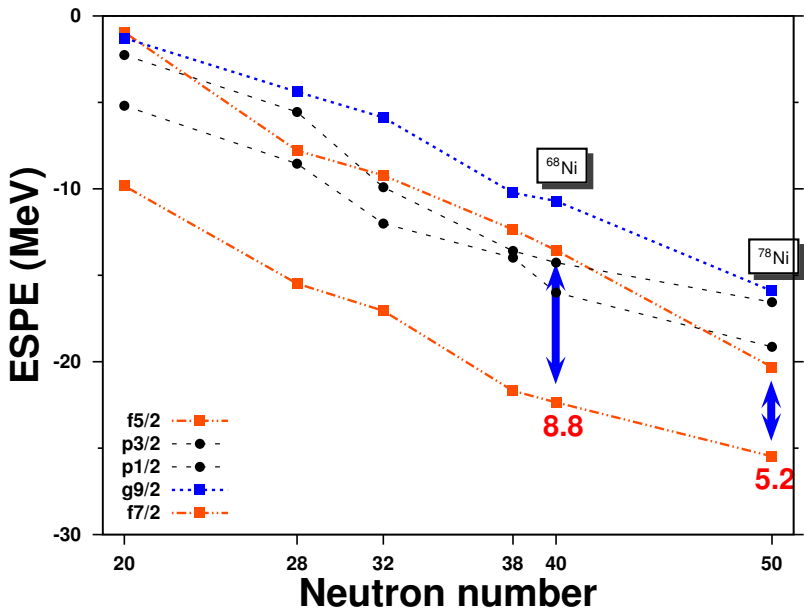
## Silicium chain



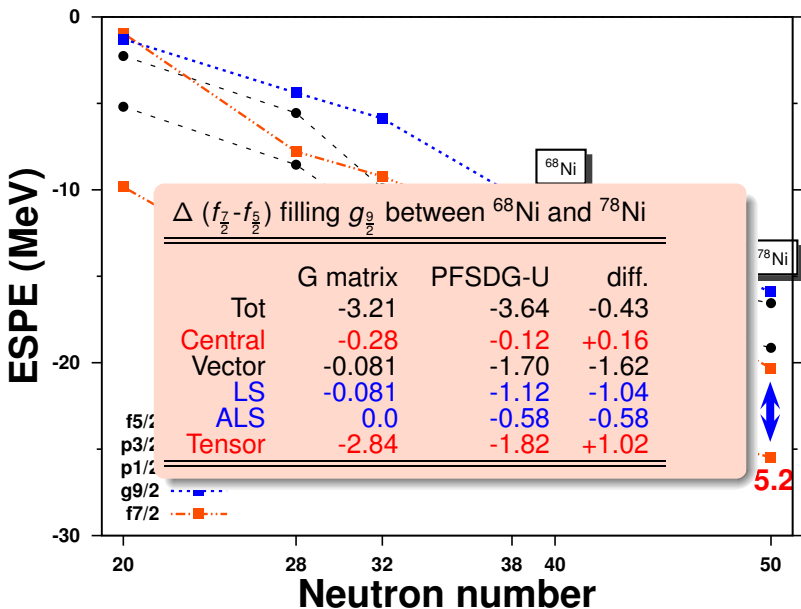
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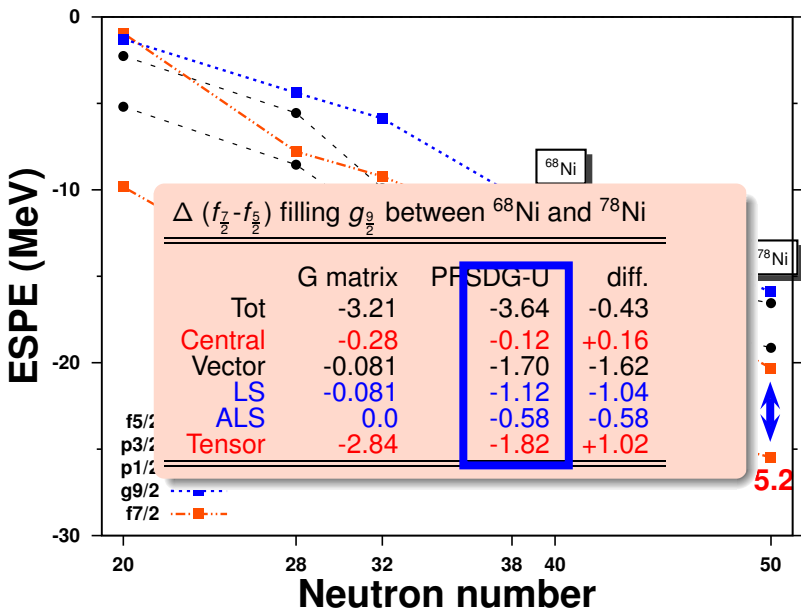
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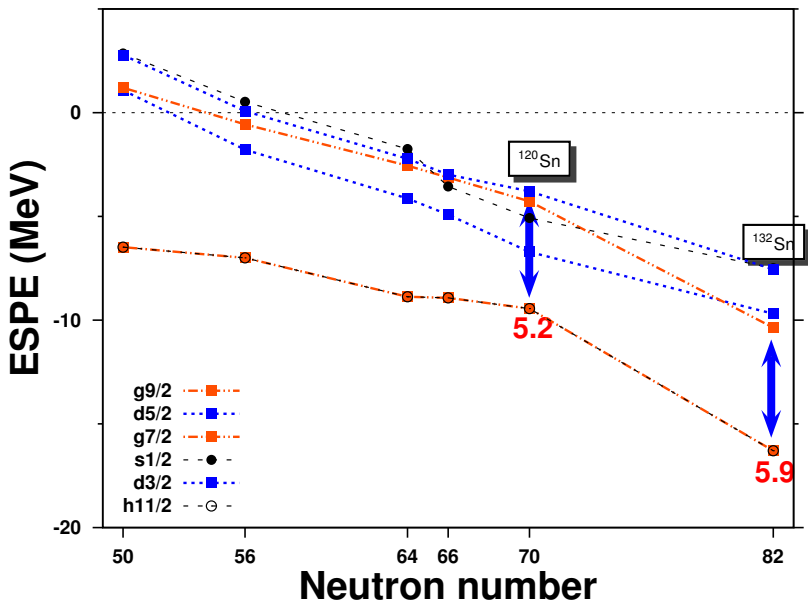
# Effective Single Particle Energies: Trends



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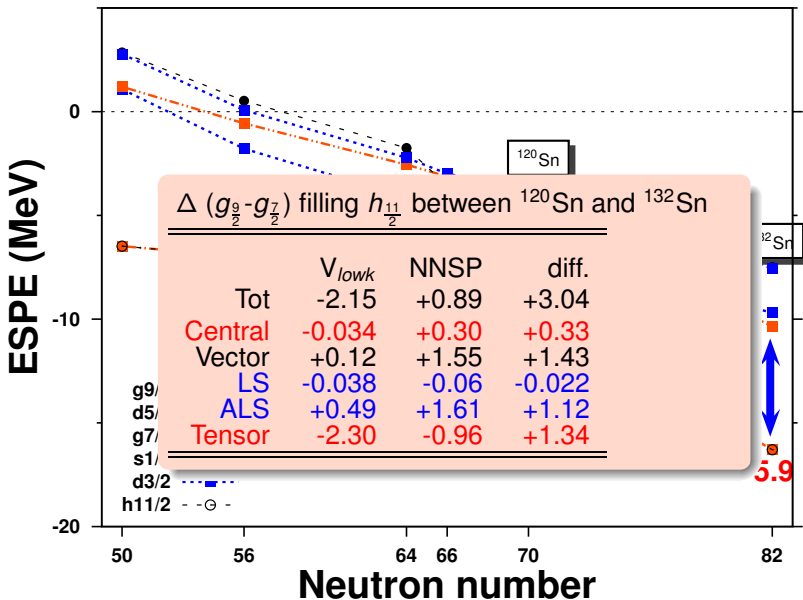


# Effective Single Particle Energies: Trends

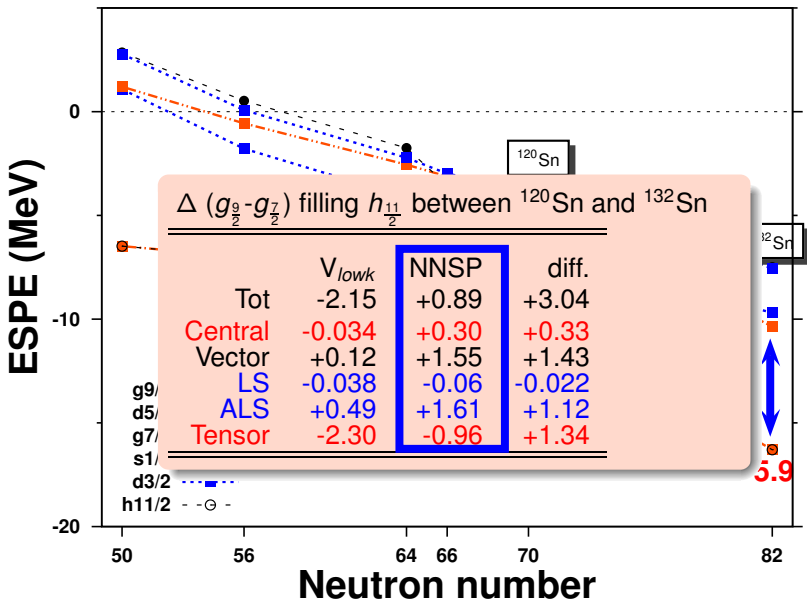




# Effective Single Particle Energies: Trends



# Effective Single Particle Energies: Trends



# Summary

- Monopole drift develops in all regions but the Interplay between correlations (pairing + quadrupole) and spherical mean-field (monopole field) determines the physics.  
It can vary from :
  - island of inversion at N=20 and N=40
  - deformation at Z=14, N=28 for  $^{42}\text{Si}$  and shell weakening at Z=28, N=50 for  $^{78}\text{Ni}$
  - deformation extending from N=40 to N=50 for Z=24-26 for  $^{74}\text{Cr}$  and  $^{76}\text{Fe}$
- The “islands of inversion” appear due to the effect of the correlations, hence they could also be called “islands of enhanced collectivity”. As quadrupole correlations are dominant in this region, most of their inhabitants are deformed rotors. Shape transitions and coexistence show up everywhere
- Quadrupole energies can be huge and understood in terms of symmetries
- Spin-Tensor Analysis show competing trends but varying significantly from light to middle mass nuclei

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